

METASTABLE STATES
IN MEDIUM - AND
HEAVY - WEIGHT NUCLEI

BY
AUGUST ALEXANDER EBEL

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METASTABLE STATES IN MEDIUM-

AND HEAVY-WEIGHT NUCLEI.

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by

August Alexander Ebel

B.A. Iowa State Teachers College
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A B S T R A C T

This thesis describes an experimental investigation of the excitation of the metastable states in In^{115} and Au^{197} by inelastic scattering of neutrons. The results are compared with those for quantum excitation of these states.

The neutron excitation curve for gold shows two pronounced discontinuities in slope at 1.11 ± 0.03 Mev and 1.14 ± 0.03 Mev, corresponding to energy levels that decay to the metastable state. The threshold for production of the isomer with neutrons is found to coincide with the energy of the metastable level.

It is shown that the spin of the metastable state must be at least $11/2$. A comparison of the cross section for the direct production of this state with the cross section for formation of the compound nucleus indicates that the state is excited by an $\ell = 3$ neutron. This fixes the spin at $11/2$ with parity opposite to that of the ground state.

An analysis of quantum-excitation data shows the 1.11 -Mev level to have a spin of $5/2$ or $7/2$ also with parity opposite that of the ground state. From the properties of the isomeric transition, the spin of the known level between the metastable level and ground is shown to have a spin of $5/2$ with parity opposite that of the ground state.

The curve obtained by neutron excitation of indium reveals a threshold at 600 ± 40 kev and other excited levels at 960 ± 40 kev and 1.37 ± 0.04 Mev. The last two have been reported from quantum excitation at 1.04 and 1.42 Mev.

TABLE I

This table contains the calculated values of the χ^2 for the various cases of the χ^2 test. The values are given for the various cases of the χ^2 test. The values are given for the various cases of the χ^2 test.

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Spin assignments of $1/2$ for the metastable state and $3/2$ or $5/2$ for the 600-kev level have been made. The parity of both levels is opposite to that of the ground state.

Confirmation is given to the existence of a beta-decay of the metastable state of In^{115} .

1. The first step is to identify the problem or question that needs to be answered. This involves understanding the context and the specific requirements of the task.

I. INTRODUCTION

ORIGIN OF THE PROBLEM

Of special interest in nuclear physics is the determination of the energy, spin, and parity of excited levels in nuclei. With only a few exceptions, the lifetimes of these levels are too short to permit a direct measurement. Instead, their characteristics must be inferred from transitions to and from the ground state. Such studies involve nuclear reactions which leave the nucleus excited. The inelastic scattering of neutrons is one such reaction. In addition to its value in the theory of nuclear structure, inelastic scattering is of considerable practical importance in the problem of nuclear shielding.

Measurements of neutron inelastic scattering are difficult because of the small cross section and the inherently large background. In a few cases, however, the lifetime of the excited state is long enough to permit its measurement to be made after the beam of incident neutrons has been removed. Thus, neutron excitation of metastable states affords a very convenient method of investigating inelastic scattering in certain nuclei.

The purpose of this thesis is to investigate the inelastic scattering levels in Au^{197} and In^{115} which decay to the metastable states of these isotopes. From an analysis of the neutron excitation curves, the spin and relative parity of many of these levels will be determined.

It should be noted that the above is a summary of the results of the study and does not represent the views of the author. The author is not responsible for the accuracy of the information contained in this document.

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HISTORY OF ISOMERISM

An isomeric state is generally defined to be an excited state of a nucleus which exists for a measurable length of time. It differs from other excited states only in the probability of decay to a lower state. With the recent advances in techniques of measurement, many states have been reclassified isomeric, and many more are being found.

The first case of isomerism was the discovery by Mahn¹ in 1921 that Pa^{234} had two half-lives, one of 1.15 minutes and the other of 6.7 hours, associated with two beta-decay energies. He was able to show that both states grew out of the beta-decay of Th^{234} . In spite of the search for other examples, it was not until the year after the discovery of artificial radioactivity in 1934 that the next isomeric pair was found. Neutron, especially slow neutron, bombardment of bromine yielded three different radioactive decay constants, although bromine was known to have only two isotopes, Br^{79} and Br^{81} . It was shown that these three periods were due to capture of neutrons giving Br^{80} , Br^{82} , and an isomer of the former.

Table I lists the lifetime and excitation energy of 23 known isomers of stable nuclei, together with the relative isotopic abundance of the stable state. Also included are the properties of the radiation arising from neutron capture which might interfere seriously with measurements of the neutron activation of the isomeric state.

With the existence of isomerism well established, the next step was a theoretical explanation by Weizsacker⁴. He established that the assumption of an excited level of a few hundred kev above the ground

1. The first step in the process of developing a business plan is to conduct a market analysis. This involves researching the industry, identifying potential customers, and understanding the competitive landscape. A thorough market analysis provides valuable insights into the viability of the business idea and helps to shape the overall strategy.

1. The first step in the process of identifying a problem is to define the problem. This involves identifying the symptoms of the problem and determining the scope of the problem. Once the problem has been defined, the next step is to identify the causes of the problem. This involves identifying the factors that are contributing to the problem and determining the underlying causes. Once the causes have been identified, the next step is to develop a plan of action. This involves identifying the steps that need to be taken to solve the problem and determining the resources that will be needed to implement the plan. Once a plan of action has been developed, the next step is to implement the plan. This involves carrying out the steps that have been identified in the plan and monitoring the progress of the implementation. Finally, the last step in the process is to evaluate the results of the implementation. This involves determining whether the problem has been solved and whether the resources have been used effectively.

1. The first step is to identify the problem or question that needs to be answered. This involves understanding the context and the specific information required.

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1. The first step is to identify the problem or issue that needs to be addressed. This involves gathering information and understanding the context of the situation.

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TABLE I

Isomers of Stable Nuclei

Isomeric Activity			Capture Activity		
Isotope	Half-Life	Excitation	Isotope	Half-Life	Rel. Abund.
		Energy MeV			
$^{72}_{32}\text{Ge}$	5×10^{-7} s	0.70			
$^{77}_{31}\text{Ga}$	17.5 s	0.15	$^{83}_{31}\text{Ga}$	67 s	3.4
$^{83}_{36}\text{Kr}$	113 m	(0.029) (0.016)	$^{84}_{36}\text{Kr}$	4.5 m	1.0
			$^{86}_{36}\text{Kr}$	74 m	4.
$^{87}_{36}\text{Sr}$	2.7 m	0.386	$^{91}_{36}\text{Sr}$	70 m	0.170
$^{93}_{41}\text{Nb}$	42 d				
$^{103}_{45}\text{Rh}$	15-43 m	0.06			
$^{107}_{47}\text{Ag}$	40 s	0.093	$^{107}_{47}\text{Ag}$	2.3 m	2.8
					52.5

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21/12/2023	17.00	0.000	12.00	21/12/2023	17.00	0.000	12.00
24/12/2023	17.00	0.000	12.00	24/12/2023	17.00	0.000	12.00
27/12/2023	17.00	0.000	12.00	27/12/2023	17.00	0.000	12.00
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$^{109}_{47}\text{Ag}$	39 d	0.008	43.7	$^{109}_{47}\text{Ag}$	24 s	2.6	47.5
$^{111}_{48}\text{Cd}$	48.7 h	0.119	12.75				
$^{113}_{48}\text{Cd}$	2.3 h		12.3				
$^{113}_{49}\text{In}$	105 h	0.393	4.2				
$^{115}_{49}\text{In}$	4.5 h	0.333	95.0	$^{115}_{49}\text{In}$	54 h	0.85	95.8
$^{119}_{50}\text{Sn}$	13 d	0.250	8.6	$^{122}_{50}\text{Sn}$	10 d	2.6	6.11
$^{125}_{52}\text{Te}$	60 d	0.125	7.0	$^{125}_{52}\text{Te}$	90 d	0.76	7.0
$^{135}_{56}\text{Ba}$	28.7 y	0.29	6.59	$^{128}_{52}\text{Te}$	32 d	1.80	31.7
$^{137}_{56}\text{Ba}$	158 s	0.663	11.32				
$^{169}_{69}\text{Tm}$	1×10^{-6} s	0.19	100				

¹⁷⁷ Hf ¹⁷⁹	19 s	0.20	(177-19.5 179-13.8)	
¹⁸¹ Ta	2×10^5 s	0.133-0.478	100	
¹⁸³ Hf	5.5 s	0.100	14.24	
¹⁸⁷ Re	0.65×10^{-6} s	0.135	62.9	
¹⁹⁷ Ir	7.4 s	0.250	100	
²⁰¹ Pb	65-68 h	1.1	1.5	²⁰¹ Pb ²⁰¹
				3.32 h 0.70
				52.3

The data in this table are taken from tabulations by Segre and Nichols² and by Seaborg and Perlman³. The capture activity, which might confuse measurements on the isomer, is listed according to the stable isotope which captures the neutron.

[illegible]

state with a spin difference several units was sufficient to predict such lifetimes.

It had long been suspected that an isomer is merely a metastable level, but the influence of atomic spectra in which large spin changes do not appear was carried over into nuclear considerations. Flugge⁵ showed this difference very clearly when he compared the atomic case having a Coulomb potential with a nuclear case assuming a potential well in a Fermi gas model. According to the build-up principle for electrons, the principal quantum numbers, n , and the azimuthal quantum numbers, ℓ , are such that, although n gets very large, nowhere does $\Delta\ell$ exceed three. In Flugge's hypothetical nuclear model, on the other hand, it is found that $\Delta\ell$ may be four or five for protons and as high as six or seven for neutrons. This scheme also shows that isomerism is unlikely in light nuclei.

More recently Mayer⁶ has proposed a model based on strong spin-orbit coupling which predicts the existence of large spin differences in low-lying levels. This explains in a rough way the "islands" of isomerism that have been observed in tabulations of isotopes. In many cases, however, the predicted spins and parity of the excited state do not agree with the experimentally determined values⁷.

EXCITATION OF METASTABLE STATES

The excitation of metastable states may be accomplished in general by two methods:

1. Electromagnetic excitations by high-energy x-rays or the electric field of charged particles; and,

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The results show that women's health status is related to their social support network. The findings suggest that women who have a strong social support network are more likely to report better health status than those who do not.

1. The Commission is of the opinion that the Commission should be authorized to conduct such investigations as may be necessary to determine the extent of the problem and to develop a plan of action to deal with it.

2. Nuclear reactions.

With the exception of a few measurements to indicate the threshold for neutron excitation, all determinations of the levels decaying to the metastable state have been made by the first method.

Excitation curves giving the energies of levels decaying to the isomer have been determined by x-rays for In^{115} , Au^{197} , Rh^{103} , Cd^{111} , and Ag. Indium is described in Chapter V and gold in Chapter IV.

The other three isomers have been studied by Wiedenbeck^{8,9}. Unfortunately, he found it necessary to use a thick target to produce x-rays, so that the decision on the existence and location of the levels is subject to large error. Figure 1 shows the level diagrams that have been reported for these nuclei.

The two stable isotopes of silver, Ag^{107} and Ag^{109} , have metastable states with very similar properties, the first having a half-life of 40 seconds and an energy of 93 kev compared to 39 seconds and 83 kev for the second. It was therefore impossible to separate the levels belonging to one isotope from the other.

Neutron excitation by the (n,n) reaction to obtain an excitation curve has been reported for In^{115} by Cohen¹⁰ in the region above 2 Mev and at four points between 600 and 1500 kev by Taschek¹¹. Neither of these investigations shows the presence of levels other than the threshold level. Wiedenbeck¹² has studied Au^{197} in the vicinity of the photon threshold. However, no careful examination has been made of the neutron excitation function of any metastable state in the region from its threshold to 2 Mev with an attempt to locate the various different levels which decay to such a state.

[illegible]

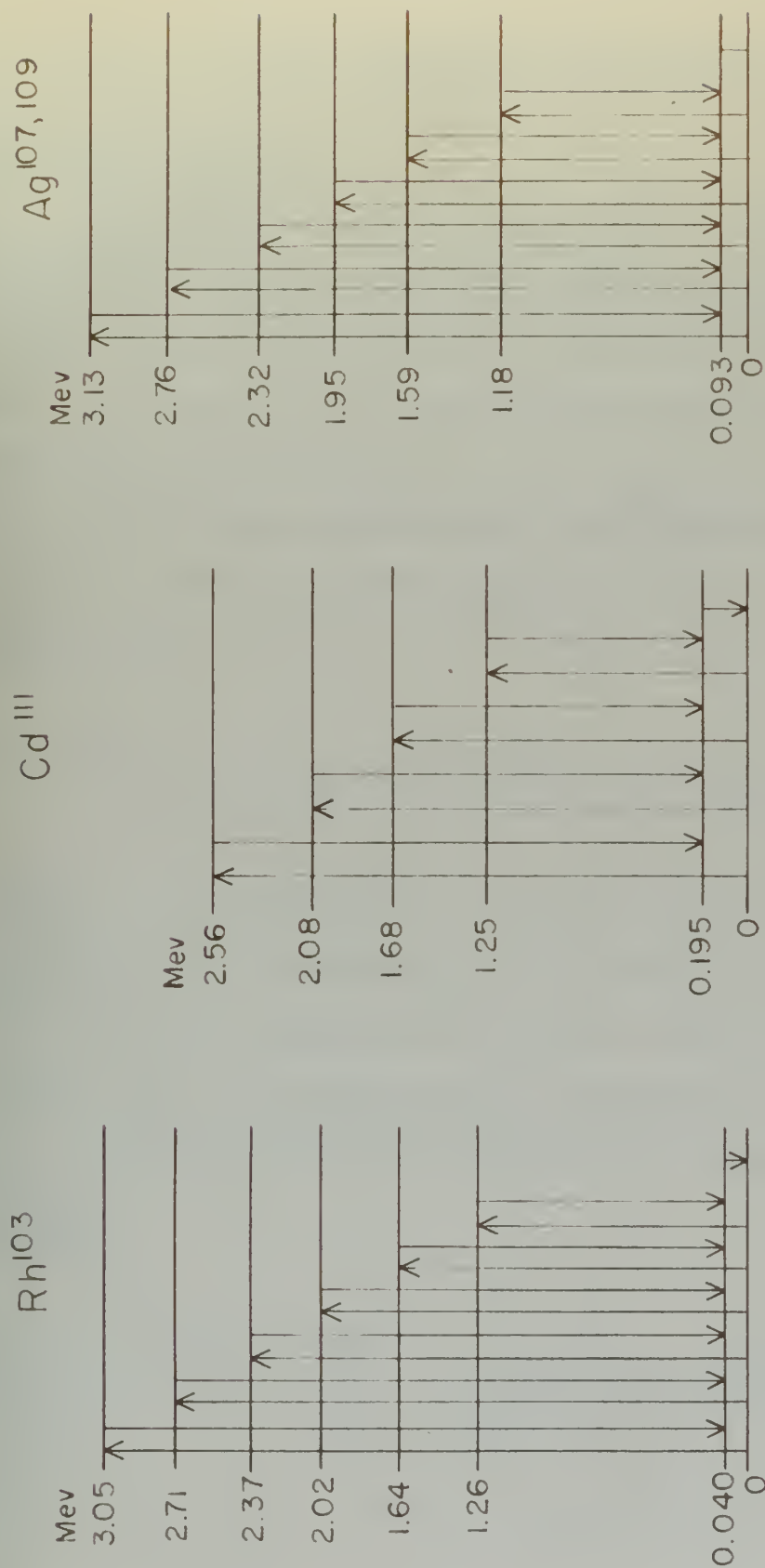


Figure 1

ENERGY LEVEL DIAGRAMS for Rh¹⁰³, Cd¹¹¹ and Ag^{107,109} PROPOSED by
 WIEDENBECK from QUANTUM EXCITATION CURVES.
 (From Phys. Rev. 67, 267 (1945) and Phys. Rev. 67, 92 (1945)).

II. THEORETICAL CONSIDERATIONS

WEISSACKER HYPOTHESIS

A quantitative explanation of the lifetime of metastable states is complicated by the lack of detailed knowledge of the nucleus. As a result, there are many expressions for the calculation of the half-life of an isomeric state, depending on the assumption of the motion of charges in the nucleus.

From classical electrodynamics, considering an oscillating system of charges, the electric field at a point R, Ω may be expressed as¹³

$$\underline{E} = -\frac{i\omega}{c^2} \frac{e^{i\omega(t-R/c)}}{R} \int \underline{j}_\perp e^{i\mathbf{k} \cdot \underline{n} \cdot \underline{r}} d\tau$$

where \underline{j}_\perp is the projection of the current vector on the plane perpendicular to the unit vector, \underline{n} , in the direction of the point R, Ω , and r, θ are the coordinates of the charge producing the current. The integration is to be carried over the distribution of charges.

According to a well-known expansion, the exponential may be expressed as

$$e^{i\mathbf{k} \cdot \underline{n} \cdot \underline{r}} = \sum_{\ell=0}^{\infty} \frac{i^\ell (k r)^\ell}{1 \cdot 3 \cdot 5 \dots (2\ell-1)} P_\ell(\cos \theta)$$

THE THERMAL EQUILIBRIUM

CHAPTER IV

A quantitative expression of the law of conservation of energy is contained in the law of thermal equilibrium of the system. In a closed system, there can be no exchange of energy with the surroundings. If the system is in thermal equilibrium with its surroundings, the temperature of the system is equal to the temperature of the surroundings.

Let us consider a system of particles, each having a certain energy. The total energy of the system is the sum of the energies of the individual particles. If the system is in thermal equilibrium with its surroundings, the total energy of the system is equal to the total energy of the surroundings.

$$E = \sum_{i=1}^N \epsilon_i = \sum_{i=1}^N \frac{1}{2} m v_i^2 = \frac{1}{2} m N \overline{v^2}$$

where ϵ_i is the energy of the i th particle, m is the mass of the particle, v_i is the velocity of the particle, and $\overline{v^2}$ is the mean square velocity of the particles. The total energy of the system is equal to the total energy of the surroundings. If the system is in thermal equilibrium with its surroundings, the total energy of the system is equal to the total energy of the surroundings.

CHAPTER V

$$E = \sum_{i=1}^N \epsilon_i = \sum_{i=1}^N \frac{1}{2} m v_i^2 = \frac{1}{2} m N \overline{v^2}$$

in the case in which the wave length of the radiation is long compared with the size of the charge distribution.

This leads to an expression for \underline{E} as the sum of integrals of various powers of $kr = r/\lambda$ corresponding to the electric field of different 2^ℓ multipole orders. The first few terms are given below:

$$\underline{E} = -\frac{1}{c^2 R} \omega e^{i\omega(t - R/c)} \left\{ \int \underline{j}_\perp P_0(\theta) d\tau \right. \\ \left. + \int \underline{j}_\perp kr P_1(\theta) d\tau + \int \underline{j}_\perp \frac{(kr)^2}{3} P_2(\theta) d\tau + \dots \right\}$$

Consideration of the magnetic field leads to a similar result.

The transition probability may be computed by forming the Poynting vector and dividing by the energy of the radiation. Collecting the terms of like powers of kr gives an expansion for W_γ of the form:

$$W_\gamma = W_{E_1} + (W_{M_1} + W_{E_2}) + (W_{M_2} + W_{E_3}) + \dots$$

in which

$$W_{M_{\ell-1}} = W_{E_\ell} = (1/\lambda)^{2\ell+1} \frac{e^2}{2} \eta^2 \frac{(r/\lambda)^{2\ell}}{[1.3.5 \dots (2\ell-1)]^2}$$

in the case of the first two cases, the results are as follows:

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The approximation to the probability for gamma-decay of the isomeric state given above was based on Bethe's¹⁴ assumption of a uniform-charge distribution throughout the nucleus. Various other models have been assumed for this derivation. Hebb and Uhlenbeck¹⁵ have assumed the radiation to be due to the electromagnetic field of a single alpha-particle moving in a potential well formed by the nucleus. Their result shows approximately the same dependence on the multipole order but includes a dependence on the charge of the nucleus. They give for the transition probability of the 2^ℓ electric or $2^{\ell-1}$ magnetic radiation:

$$\lambda = 8\left(\frac{1}{\kappa}\right)\left(\frac{\ell+1}{\ell}\right) \frac{e^2}{\hbar} \left(\frac{(r/\kappa)^{2\ell}}{2(1.3.5\dots(2\ell+1)^2)} \right)$$

Flügge¹⁶ does not consider the charge distribution performing a linear oscillation, but instead he pictures the radiation arising from a rotational excitation of a charged drop. This leads to another expression similar to others but differing in the constants arising from the nuclear properties and is of the form:

$$\lambda = Z^2 \frac{e^2}{\hbar c} \frac{\hbar}{\Theta} \frac{E}{Mc^2} (r/\kappa)^{2\ell}$$

where Θ is the moment of inertia and M the rest mass of the nucleus.

The accuracy of any of these formulas is sufficient to determine the ℓ value of the transition in most cases because of the vast dif-

The investigation of the conditions for the existence of the solutions of the system (1) is carried out in the following way. First, we shall find the conditions for the existence of the solutions of the system (1) in the case when the functions $f_i(x)$ are continuous and the functions $g_i(x)$ are piecewise continuous. Then, we shall find the conditions for the existence of the solutions of the system (1) in the case when the functions $f_i(x)$ are piecewise continuous and the functions $g_i(x)$ are continuous. Finally, we shall find the conditions for the existence of the solutions of the system (1) in the case when the functions $f_i(x)$ and $g_i(x)$ are both piecewise continuous.

$$x = \frac{1}{2} \left(\frac{1}{2} + \frac{1}{2} \right) = 1$$

It is not difficult to see that the conditions for the existence of the solutions of the system (1) are satisfied in the case when the functions $f_i(x)$ are continuous and the functions $g_i(x)$ are piecewise continuous. This is because the functions $f_i(x)$ are continuous and the functions $g_i(x)$ are piecewise continuous, and the conditions for the existence of the solutions of the system (1) are satisfied in the case when the functions $f_i(x)$ are continuous and the functions $g_i(x)$ are piecewise continuous.

$$x = \frac{1}{2} \left(\frac{1}{2} + \frac{1}{2} \right) = 1$$

It is not difficult to see that the conditions for the existence of the solutions of the system (1) are satisfied in the case when the functions $f_i(x)$ are continuous and the functions $g_i(x)$ are piecewise continuous. This is because the functions $f_i(x)$ are continuous and the functions $g_i(x)$ are piecewise continuous, and the conditions for the existence of the solutions of the system (1) are satisfied in the case when the functions $f_i(x)$ are continuous and the functions $g_i(x)$ are piecewise continuous.

ference in lifetimes for different multipole-order transitions. However, an error of as much as 10^ℓ to $10^{2\ell}$ may be expected due to the inaccuracy of the particular model. None of these models appear at present to have much advantage over the others.

ANGULAR MOMENTUM SELECTION RULE

By a consideration of the fields of the region in which they do not decay as R^{-1} close to the source, Heitler¹⁷ has shown that quantization of these fields results in the radiation of a 2^ℓ electric or magnetic pole having an angular momentum of $\ell \hbar$ with respect to the multipole origin.

Since vector angular momentum of the nucleus must be conserved, the emission of radiation changes the angular momentum in accordance with the multipole order of the radiation. If the initial and final states of the nucleus differ by more than one unit of momentum in absolute magnitude, this places a lower limit on the order of the multipole. Conversely, if the lifetime of the excited state indicates that the lowest multipole order present is 2^ℓ , it can be stated that $|I| - |I'| = \ell$.

Although 2^ℓ may be the lowest-order multipole radiation possible, higher-order radiation is of course also possible up to the limit determined by $|I| + |I'|$. However, barring a coincidence of charge distribution that reduces the strength of the 2^ℓ pole far below that which would be predicted from the size of the nucleus, only the lowest-order radiation would be observed because of its much greater probability.

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PARITY SELECTION RULE

Another selection rule that must always be obeyed is the conservation of parity in a transition. It can be shown by a consideration of the symmetry properties of the eigenfunction of the nucleus that emission of a $2^{\ell-1}$ magnetic or 2^{ℓ} electric radiation always involves a parity change for the nucleus if ℓ is odd and never if ℓ is even. Thus, although these transitions involve the same probability in the rough derivation above, it is impossible for both to occur between the same two states.

INTERNAL CONVERSION

Any comparison of these results with theory involves a correction for the competing reaction, internal conversion. Taylor and Mott¹⁸ have shown that internal conversion, being a one-step process and not a photoelectric effect of the outgoing gamma-ray on the atomic electrons, is in fact a competing reaction reducing the lifetime of the excited nucleus. Therefore, it is correct to define the internal-conversion coefficient as the ratio of the conversion electrons to the number of gamma-ray quanta and not as the ratio to their sum.

The calculation of these coefficients involves an evaluation of the probability of transmitting to the atomic electrons the excitation energy of the nucleus by a perturbation of their wave functions, this perturbation being caused by a multipole in the nucleus of such strength that it radiates one quantum per second. Since the fields due to the higher multipole orders decrease much more rapidly with distance than do

those of the lower orders, this factor provides another check on the order of an observed transition.

Such a calculation in the general case is so complicated that the formulae derived are very restricted in their application. Dancoff and Morrison¹⁹ have obtained expressions for the case in which the emitted electrons can be treated nonrelativistically and in the relativistic case in which the binding energy of the electrons in the shell can be neglected. However, numerical integration of the wave functions by Rose²⁰ shows that, even in this limited region, the predicted values may be in error by a factor of three or more especially for the higher-order multipoles and magnetic conversion. There are several other approximate formulae available, but these are no better when checked with Rose's data. Therefore, at the present time, all calculations of internal conversion coefficients must involve numerical integration of the Schrodinger equation.

The accurate experimental determination of the internal conversion coefficient is difficult because of the uncertainty of the relative efficiencies of the detectors for electrons and gamma-rays. The ratio of the coefficients for the K and L shells is much more accurately determined and is a more sensitive function of the transition, particularly in separating electric and magnetic transitions of similar probabilities. For example, gold has a measured K to L ratio of 3.4. Assuming electric 2^4 transition gives 0.51 for this factor, and a magnetic 2^3 pole has 4.8. The transition is therefore considered to be magnetic.

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1994. *Diary of a mad biologist: a collection of madcap tales*.

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1. The first step is to identify the problem or question that needs to be answered. This involves understanding the context and the specific requirements of the task.

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TABLE 1. *Continued*

Axel and Dancoff²¹ have compared the lifetimes of 58 known isomers with the theoretical lifetimes predicted by the Weissacker hypothesis corrected for internal conversion and find that these isomers do in fact fall into groups characteristic of $\ell = 4$ and $\ell = 5$ spin change, with one case, that of Pb^{204} , corresponding to $\ell = 6$. This agreement indicates that the hypothesis is valid.

NEUTRON EXCITATION OF METASTABLE LEVELS

The neutron excitation of a metastable state is a special case of inelastic scattering of neutrons. The process is shown diagrammatically in Figure 2. On the left is shown the bombarded nucleus with mass number A , and on the right is the compound nucleus with mass number $A + 1$ more tightly bound by the binding energy of the additional particle (on the order of 8 Mev). Nucleus A has a ground state G and a first excited level E .

Bombardment by a neutron with energy less than the energy of level E is shown by the dotted line. In this case, the compound state can emit a neutron leaving the nucleus A only in the ground state. However, if the bombarding neutron has a higher energy than does the excited level, the compound nucleus can decay to level E or to the ground state, as shown by the solid lines. The lowest neutron energy that excites level E is the energy of that level, provided the neutron is able to carry the difference in angular momentum between states G and E . If this is not the case, it may be necessary to excite a higher state E' , whose angular momentum lies between that of E and G . Level E may then be excited by the emission of a gamma-ray by level E' .

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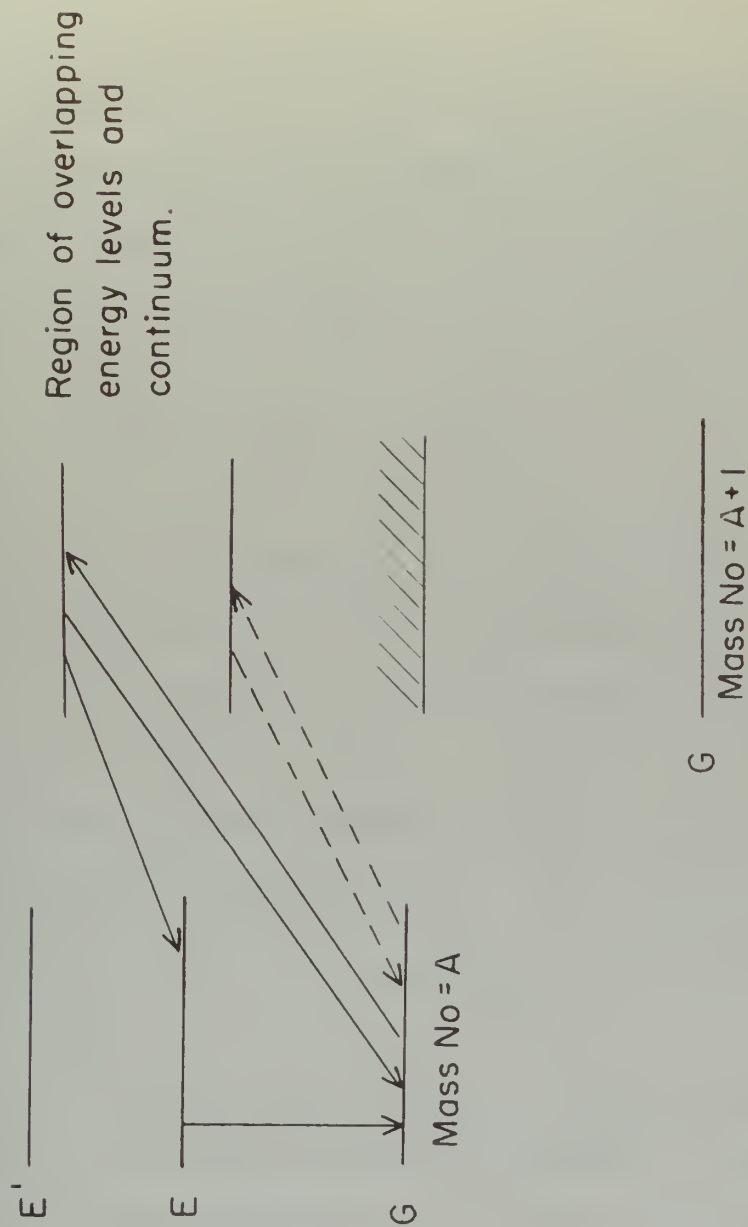


Figure 2
SCHEMATIC DIAGRAM of NEUTRON SCATTERING with NEUTRON
ENERGY LESS (broken lines) and GREATER (solid lines) THAN
FIRST EXCITED LEVEL E .

The cross section for exciting level E is the product of the cross section of exciting the compound nucleus, σ_c , and the probability, W_E , of that nucleus decaying to level E .

Blatt and Weisskopf²² have derived an expression for the cross section for formation of the compound nucleus as a function of the energy and orbital momentum of the neutron with respect to the target nucleus. This is given to be:

$$\sigma_c(\ell) = (2\ell + 1) \pi \lambda^2 T_\ell$$

where λ is the wave length of the incident neutron and T_ℓ is the transmission coefficient for the neutron through the "centrifugal" potential and also includes the reflection due to the energy discontinuity at the nuclear surface.

Assuming a square potential well for the nucleus, the transmission coefficient is shown to be:

$$T_\ell \approx 4 \frac{x}{X} v_\ell$$

for the case in which the wave number K of the neutron in the nucleus is large compared with the wave number k outside the nucleus. In this expression $x = kR$ and $X = KR$, where R is the nuclear radius and v_ℓ is a factor determined by the wave function of the neutron beam at the surface of the nucleus.

The values of v_ℓ may be calculated from a tabulation of Blatt and Weisskopf giving:

The same condition for finding level 1 in the context of the
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 The same condition for finding level 1 in the context of the
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$$x^2 = (x+1)^2 - 2x - 1$$

where x is the new length of the standard deviation and 1 is the
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 deviation and 1 is the new length of the standard deviation.
 Assuming a square potential well for the motion, the potential
 energy coefficient is given by

$$V \approx \frac{x}{L} \approx \frac{x}{L}$$

for the case in which the new length L of the motion is the
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 the motion. In this case, $x = 1$ and $L = 1$, so $V \approx 1$ is the
 new length of the motion and V is a linear coefficient of the new length
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 The value of V is given by a relation of the
 new length of the motion.

$$\begin{aligned}
v_0 &= 1 \\
v_1 &= x^2/(1+x^2) \\
v_2 &= x^4/(9+3x^2+x^4) \\
v_3 &= x^6/(225+15x^2+(x^4+x^6)) \\
v_\ell &\rightarrow x^{2\ell}/[1.3.5...(2\ell-1)]^2
\end{aligned}$$

The last expression gives the value of v_ℓ in the limit $x \ll \ell$.

It should be noted that the angular momentum a neutron can carry into the compound nucleus for a given energy is much greater than would be predicted from a simple quantization of the impact parameter. An 800-kev neutron incident on Au^{197} can react only when it has orbital angular momentum $\ell = 0$ or $\ell = 1$, according to the naive approach. Actually, the cross section for $\ell = 2$ is greater than for $\ell = 0$ and over half that for $\ell = 1$. $\ell = 3$ is more than one-tenth and even $\ell = 4$ is more than one-hundredth of the $\ell = 1$ value.

The probability for decay to a given state from the compound nucleus follows from a consideration of the reciprocity theorem which states that

$$\sigma(A \rightarrow B)/\lambda_A^2 = \sigma(B \rightarrow A)/\lambda_B^2$$

connects the cross section for a transition from state A to state B with the reverse transition. Neglecting all methods of decay of the compound nucleus other than neutron emission, the fraction of the compound nuclei which decays to state 1 with energy E_1 above the ground state in competition with n other states 1 of energy E_1 is given by:

$$P_1 = \frac{(E - E_1) \sigma_c(1)}{\sum_{i=1}^N (E - E_i) \sigma_c(i)}$$

where $\sigma_c(i)$ is the cross section for the formation of the compound nucleus from level i and E is the energy available for the transition from the compound nucleus to the ground state of the product nucleus.

In a small energy interval just above the threshold for level "1" the sum in the denominator will change only slightly, so that in this interval the probability is proportional to $(E - E_1) \sigma_c(1)$. As the energy of the compound nucleus is increased, the increase in the sum can no longer be neglected, and the probability will be less than predicted by the proportionality.

An inspection of the form of $\sigma_c^{(\ell)}$ in the vicinity of the threshold shows that for the case $\ell = 0$ the cross section $\sigma_c^{(0)}$ decreases as $(E - E_{th})^{-1/2}$; therefore, the probability of decay to that state increases as $(E - E_{th})^{1/2}$. For $\ell = 1$, for the neutron emitted in the decay, the probability increases as $(E - E_{th})^{3/2}$ and in general as $(E - E_{th})^{(\ell + 1/2)}$. Thus, it is easily recognized if the decay to an excited state is accomplished by the emission of an $\ell = 0$ neutron because the excitation curve will approach the threshold with a vertical tangent; whereas in cases of higher ℓ , the tangent will be horizontal.

This type of calculation applies directly to the case in which the compound nucleus decays to the metastable state itself. In many

cases, however, this transition is so improbable because of spin changes involved that the compound nucleus decays to an excited state with an intermediate spin, and that excited state decays in turn to the metastable level. The cross section for production of the isomer must then be corrected considering the relative probability of decay from the excited state to the metastable state and the other available levels. For this, the calculations described earlier in this chapter according to the Weissacker hypothesis may be used.

COMPARISON OF NEUTRON AND PHOTON EXCITATION

Neutron and photon excitation of levels above the ground state differ in two fundamental respects. Since a neutron is re-emitted in neutron excitation, it is not necessary that the neutron have exactly the energy of the level, whereas, in the case of photons, the excitation is a line absorption. Hence, for photons, it is necessary to use a broad spectrum energy source, and increasing the photon energy gives an isochromat of the source. This makes the shape of the excitation curve dependent upon the photon source and not upon the excitation process.

The second difference is found in consideration of the angular momentum that can be carried in to change the nuclear spin from that in the ground state to that in the excited state. A dipole radiation can contribute only one unit and a quadrupole only two units. Excitation by higher-order multipole radiation is so unlikely that it is not detectable in the ordinary investigations. The same rules that make the transition from an excited state to the ground level by high

course, however, with provision as to the manner of its
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AMENDMENT OF SECTION 100 (1) (a) (i)

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multipole radiation improbable hold for the case of exciting that state from the ground level by high-order multipole radiation. Therefore, it is impossible to excite a metastable state of moderate lifetime directly by photons.

However, neutrons may carry in several units of angular momentum depending upon their energy and the size of the target nucleus. In the case of Au^{197} , a 600-kev neutron can change the spin of the nucleus four units and is able to excite the metastable state directly. On the other hand, neutrons can also carry in small angular momentum changes with a high degree of probability, so that any state which can be excited by photons can also be excited by neutrons; and, in addition, neutrons can excite many states which cannot be excited by photons.

1. The first step in the process of identifying a problem is to determine the nature of the problem. This involves a thorough understanding of the situation and the factors that are contributing to the problem. Once the nature of the problem is understood, the next step is to identify the causes of the problem. This involves a detailed analysis of the situation and the factors that are contributing to the problem. Once the causes of the problem are identified, the next step is to develop a plan of action to address the problem. This involves determining the steps that need to be taken to solve the problem and the resources that will be required to implement the plan. Finally, the last step in the process is to implement the plan and monitor the results. This involves putting the plan into action and tracking the progress of the solution to ensure that the problem is resolved.

1. The first point to be noted is that the Commission has not yet received any information from the Government of the United Kingdom regarding the proposed amendments to the Bill. It is therefore not possible to say whether or not the amendments are necessary or desirable.

III. EXPERIMENTAL EQUIPMENT

NEUTRON SOURCE

The neutrons used for activation of the metastable state were obtained from the $\text{Li}^7(\text{p},\text{n})\text{Be}^7$ reaction, using protons from the M.I.T. Rockefeller generator. This reaction was chosen because it has a comparatively large cross section and provides nearly monoenergetic neutrons..

Since it was desired to study isomeric production near the threshold for this process, it was desirable that the reaction be endoergic so that neutrons with low energies could be obtained. With a Q -value of -1.63 Mev, the $\text{Li}^7(\text{p},\text{n})$ reaction satisfies this condition.

The presence of a lower-energy neutron group with relative intensity of about 10 percent has recently been reported²³ at 384 keV below the main group. Since the excitation curve for the metastable state rises quite rapidly, this lower-energy group causes no ambiguity in the results. It is of interest that the Au^{197*} excitation curve (Figure 13) shows a slight rise within the statistical error at about 430 keV above threshold; this may be due to this effect, but the results are not conclusive.

The target was prepared by evaporating metallic lithium (92.6 percent Li^7) on a 10-mil tantalum backing. To prevent the proton beam from boiling the lithium off, the target was rotated with the beam striking approximately one inch off center. An even coating was obtained by rotating the target while the lithium was being deposited.

The majority of the members of the committee are now
appointed from the 14th District, which includes the 1st, 2nd,
3rd, 4th, 5th, 6th, 7th, 8th, 9th, 10th, 11th, 12th, 13th, 14th, 15th, 16th, 17th, 18th, 19th, 20th, 21st, 22nd, 23rd, 24th, 25th, 26th, 27th, 28th, 29th, 30th, 31st, 32nd, 33rd, 34th, 35th, 36th, 37th, 38th, 39th, 40th, 41st, 42nd, 43rd, 44th, 45th, 46th, 47th, 48th, 49th, 50th, 51st, 52nd, 53rd, 54th, 55th, 56th, 57th, 58th, 59th, 60th, 61st, 62nd, 63rd, 64th, 65th, 66th, 67th, 68th, 69th, 70th, 71st, 72nd, 73rd, 74th, 75th, 76th, 77th, 78th, 79th, 80th, 81st, 82nd, 83rd, 84th, 85th, 86th, 87th, 88th, 89th, 90th, 91st, 92nd, 93rd, 94th, 95th, 96th, 97th, 98th, 99th, 100th, 101st, 102nd, 103rd, 104th, 105th, 106th, 107th, 108th, 109th, 110th, 111th, 112th, 113th, 114th, 115th, 116th, 117th, 118th, 119th, 120th, 121st, 122nd, 123rd, 124th, 125th, 126th, 127th, 128th, 129th, 130th, 131st, 132nd, 133rd, 134th, 135th, 136th, 137th, 138th, 139th, 140th, 141st, 142nd, 143rd, 144th, 145th, 146th, 147th, 148th, 149th, 150th, 151st, 152nd, 153rd, 154th, 155th, 156th, 157th, 158th, 159th, 160th, 161st, 162nd, 163rd, 164th, 165th, 166th, 167th, 168th, 169th, 170th, 171st, 172nd, 173rd, 174th, 175th, 176th, 177th, 178th, 179th, 180th, 181st, 182nd, 183rd, 184th, 185th, 186th, 187th, 188th, 189th, 190th, 191st, 192nd, 193rd, 194th, 195th, 196th, 197th, 198th, 199th, 200th, 201st, 202nd, 203rd, 204th, 205th, 206th, 207th, 208th, 209th, 210th, 211th, 212th, 213th, 214th, 215th, 216th, 217th, 218th, 219th, 220th, 221st, 222nd, 223rd, 224th, 225th, 226th, 227th, 228th, 229th, 230th, 231st, 232nd, 233rd, 234th, 235th, 236th, 237th, 238th, 239th, 240th, 241st, 242nd, 243rd, 244th, 245th, 246th, 247th, 248th, 249th, 250th, 251st, 252nd, 253rd, 254th, 255th, 256th, 257th, 258th, 259th, 260th, 261st, 262nd, 263rd, 264th, 265th, 266th, 267th, 268th, 269th, 270th, 271st, 272nd, 273rd, 274th, 275th, 276th, 277th, 278th, 279th, 280th, 281st, 282nd, 283rd, 284th, 285th, 286th, 287th, 288th, 289th, 290th, 291st, 292nd, 293rd, 294th, 295th, 296th, 297th, 298th, 299th, 300th, 301st, 302nd, 303rd, 304th, 305th, 306th, 307th, 308th, 309th, 310th, 311th, 312th, 313th, 314th, 315th, 316th, 317th, 318th, 319th, 320th, 321st, 322nd, 323rd, 324th, 325th, 326th, 327th, 328th, 329th, 330th, 331st, 332nd, 333rd, 334th, 335th, 336th, 337th, 338th, 339th, 340th, 341st, 342nd, 343rd, 344th, 345th, 346th, 347th, 348th, 349th, 350th, 351st, 352nd, 353rd, 354th, 355th, 356th, 357th, 358th, 359th, 360th, 361st, 362nd, 363rd, 364th, 365th, 366th, 367th, 368th, 369th, 370th, 371st, 372nd, 373rd, 374th, 375th, 376th, 377th, 378th, 379th, 380th, 381st, 382nd, 383rd, 384th, 385th, 386th, 387th, 388th, 389th, 390th, 391st, 392nd, 393rd, 394th, 395th, 396th, 397th, 398th, 399th, 400th, 401st, 402nd, 403rd, 404th, 405th, 406th, 407th, 408th, 409th, 410th, 411th, 412th, 413th, 414th, 415th, 416th, 417th, 418th, 419th, 420th, 421st, 422nd, 423rd, 424th, 425th, 426th, 427th, 428th, 429th, 430th, 431st, 432nd, 433rd, 434th, 435th, 436th, 437th, 438th, 439th, 440th, 441st, 442nd, 443rd, 444th, 445th, 446th, 447th, 448th, 449th, 450th, 451st, 452nd, 453rd, 454th, 455th, 456th, 457th, 458th, 459th, 460th, 461st, 462nd, 463rd, 464th, 465th, 466th, 467th, 468th, 469th, 470th, 471st, 472nd, 473rd, 474th, 475th, 476th, 477th, 478th, 479th, 480th, 481st, 482nd, 483rd, 484th, 485th, 486th, 487th, 488th, 489th, 490th, 491st, 492nd, 493rd, 494th, 495th, 496th, 497th, 498th, 499th, 500th, 501st, 502nd, 503rd, 504th, 505th, 506th, 507th, 508th, 509th, 510th, 511th, 512th, 513th, 514th, 515th, 516th, 517th, 518th, 519th, 520th, 521st, 522nd, 523rd, 524th, 525th, 526th, 527th, 528th, 529th, 530th, 531st, 532nd, 533rd, 534th, 535th, 536th, 537th, 538th, 539th, 540th, 541st, 542nd, 543rd, 544th, 545th, 546th, 547th, 548th, 549th, 550th, 551st, 552nd, 553rd, 554th, 555th, 556th, 557th, 558th, 559th, 560th, 561st, 562nd, 563rd, 564th, 565th, 566th, 567th, 568th, 569th, 570th, 571st, 572nd, 573rd, 574th, 575th, 576th, 577th, 578th, 579th, 580th, 581st, 582nd, 583rd, 584th, 585th, 586th, 587th, 588th, 589th, 590th, 591st, 592nd, 593rd, 594th, 595th, 596th, 597th, 598th, 599th, 600th, 601st, 602nd, 603rd, 604th, 605th, 606th, 607th, 608th, 609th, 610th, 611th, 612th, 613th, 614th, 615th, 616th, 617th, 618th, 619th, 620th, 621st, 622nd, 623rd, 624th, 625th, 626th, 627th, 628th, 629th, 630th, 631st, 632nd, 633rd, 634th, 635th, 636th, 637th, 638th, 639th, 640th, 641st, 642nd, 643rd, 644th, 645th, 646th, 647th, 648th, 649th, 650th, 651st, 652nd, 653rd, 654th, 655th, 656th, 657th, 658th, 659th, 660th, 661st, 662nd, 663rd, 664th, 665th, 666th, 667th, 668th, 669th, 670th, 671st, 672nd, 673rd, 674th, 675th, 676th, 677th, 678th, 679th, 680th, 681st, 682nd, 683rd, 684th, 685th, 686th, 687th, 688th, 689th, 690th, 691st, 692nd, 693rd, 694th, 695th, 696th, 6

Since it was desired to study the effect of the amount of the
oil on the process, it was decided to use the amount of oil
as first mentioned with the exception of the amount of oil. The
oil was the same as the oil used in the first experiment.

[illegible][illegible]

Thus, the problem of installing the target without the lithium's being oxidized was eliminated.

The thickness of the targets used was of the order of 15 to 20 kev, as measured by the conventional method²⁴ of taking the thickness to be the energy between the neutron threshold and the first geometrical peak in the yield curve at 0 degrees.

ROCKEFELLER GENERATOR

The proton accelerator was an electrostatic generator originally used to accelerate electrons. During the past two years, this generator has been converted into a positive-ion accelerator by the Nuclear Shielding Group. Unfortunately, the beam emerged in a room located under ground so that the scattered neutron flux is quite high. The target was located 35 inches from the analyzing magnet and only 39 inches from the concrete floor. In other directions, the situation was considerably more favorable. For the measurements being reported, however, scattered neutrons offered no problem because the exposures were made in the region of high direct flux close to the target, and the scattered flux, being of lower energy, had a smaller cross section for interaction.

The beam was defined by crossed slits above the analyzing magnet and by horizontal slits beyond. Openings of 1 millimeter in the exit slits corresponded to an energy resolution of 0.1 percent in the proton beam. However, for the measurements on gold, it was necessary to open the slits wider to obtain sufficient neutron flux.

There are several other factors which may be considered in the selection of a site for a new plant. These are: (1) the availability of raw materials, (2) the availability of labor, (3) the availability of capital, (4) the availability of transportation, (5) the availability of water, (6) the availability of power, (7) the availability of land, (8) the availability of a good location, (9) the availability of a good climate, (10) the availability of a good market.

It is to be noted that the above is a summary of the information
available to the Bureau at the time of the investigation and is not
intended to be a final report. The Bureau is continuing its
investigation and will report the results of its findings as they
develop.

[illegible]

The film was released by Academy Award-nominee John Ford, who directed the original 1939 film. The film was released by Academy Award-nominee John Ford, who directed the original 1939 film.

The proton energies were determined by a generating voltmeter indicating on a 6 inch fan-type microammeter. This meter could be read to 0.1 of a division (about 5 kev in neutron energy), and measurements on the $\text{Li}^7(\text{p},\text{n})$ threshold indicated that it was reproducible within this amount over the period of a run. Over a period of several weeks, however, the calibration on the generating voltmeter changed by as much as several percent; therefore, it was always checked against the $\text{Li}^7(\text{p},\text{n})$ threshold (1.882 Mev) at the beginning and the end of each run.

In addition to this calibration point, the circuit has been checked for linearity using $\text{Li}^7(\text{p}_2,\text{n})$, in which the proton is accelerated in a singly charged H_2 molecule, giving a threshold at 3.764 Mev, twice the $\text{Li}^7(\text{p},\text{n})$ threshold, the $\text{H}^3(\text{p},\text{n})$ with a threshold at 1.019 Mev, and the $\text{C}^{13}(\text{p},\text{n})$ with a threshold at 3.236 Mev. It was found that the calibration of the meter against proton energy was linear in the region of interest, although the linear portion did not extrapolate to zero.

From the proton energy, the neutron energy can be calculated by application of conservation of momentum and energy. A tabulation of this calculation prepared by Willard²⁵ was used in this investigation.

LONG COUNTER

The neutron flux was measured, using a long counter of design similar to that of Hanson and McKibben²⁶. This counter, shown with dimensions in Figure 3, consists of a cylinder of paraffin shielded by cadmium surrounding a BF_3 counter. The BF_3 counter was 1 inch in diameter with an active volume 12 inches long. The filling was of 96 percent enriched BF_3 at 50 cm. of mercury pressure.

The first part of the paper is devoted to the study of the
 problem of the existence of a solution of the system
 (1.1) in the case of a general boundary value problem.
 In the second part of the paper we study the problem of the
 existence of a solution of the system (1.1) in the case of a
 boundary value problem of the second kind. In the third part
 of the paper we study the problem of the existence of a
 solution of the system (1.1) in the case of a boundary value
 problem of the third kind. In the fourth part of the paper
 we study the problem of the existence of a solution of the
 system (1.1) in the case of a boundary value problem of the
 fourth kind. In the fifth part of the paper we study the
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 the case of a boundary value problem of the fifth kind. In
 the sixth part of the paper we study the problem of the
 existence of a solution of the system (1.1) in the case of a
 boundary value problem of the sixth kind. In the seventh part
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 we study the problem of the existence of a solution of the
 system (1.1) in the case of a boundary value problem of the
 eighth kind. In the ninth part of the paper we study the
 problem of the existence of a solution of the system (1.1) in
 the case of a boundary value problem of the ninth kind. In
 the tenth part of the paper we study the problem of the
 existence of a solution of the system (1.1) in the case of a
 boundary value problem of the tenth kind. In the eleventh part
 of the paper we study the problem of the existence of a
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 problem of the eleventh kind. In the twelfth part of the
 paper we study the problem of the existence of a solution of
 the system (1.1) in the case of a boundary value problem of
 the twelfth kind. In the thirteenth part of the paper we
 study the problem of the existence of a solution of the system
 (1.1) in the case of a boundary value problem of the
 thirteenth kind. In the fourteenth part of the paper we
 study the problem of the existence of a solution of the system
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 fourteenth kind. In the fifteenth part of the paper we study
 the problem of the existence of a solution of the system (1.1)

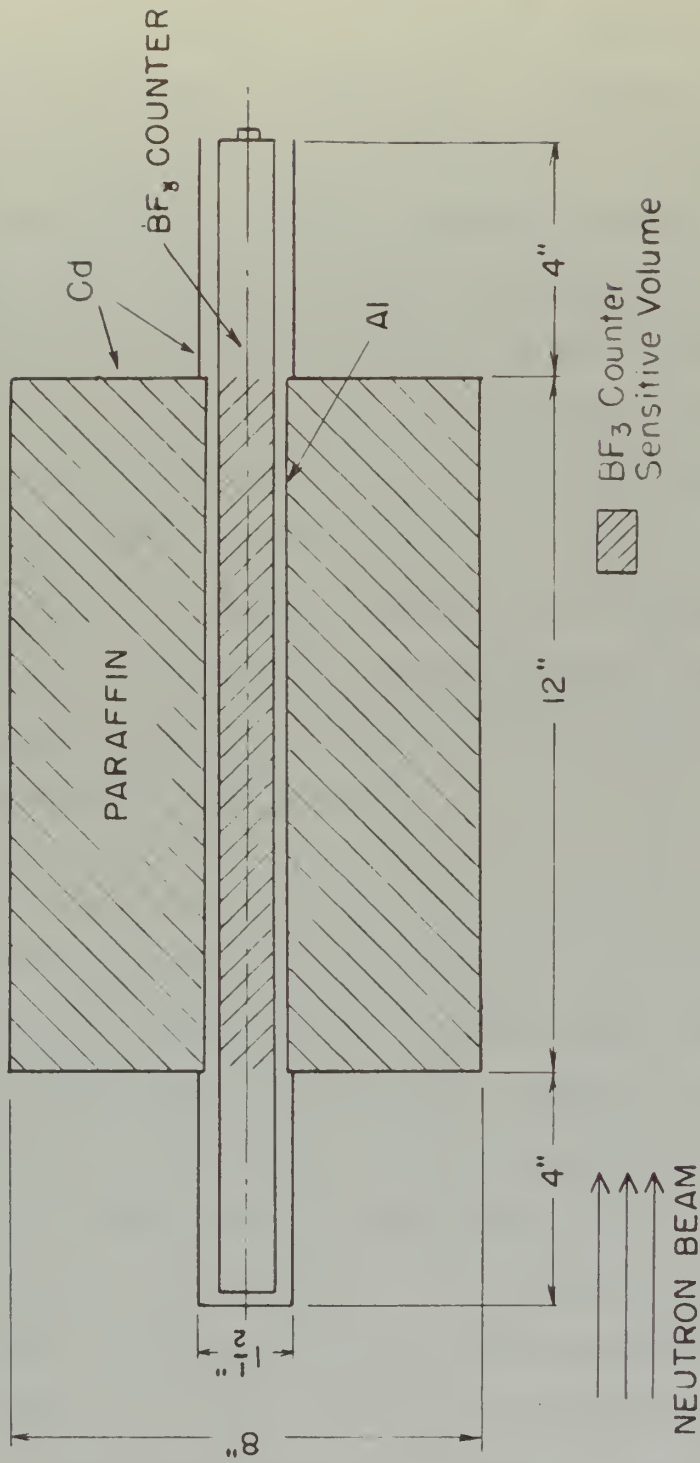


Figure 3. CROSS-SECTION of LONG COUNTER

Willard²⁵ compared the yield curve of the $\text{Li}^7(\text{p},\text{n})$ reaction as determined by this long counter with the curve published for a 10-kev target by Hanson, Taschek, and Williams²⁷. By assuming their yield curve to be correct, Willard was able to determine a calibration for the long counter. Figure 4 shows the relative response as a function of the energy of the incident neutrons. In addition, he found that, when located one meter from the target at 0 degrees, the counter gave 60 counts per microcoulomb of protons per kev target thickness in the flat portion of the $\text{Li}^7(\text{p},\text{n})$ yield curve. The target thickness was measured by the initial rise in the yield curve. This gives an absolute calibration of the counter when compared to Hanson's curve.

The associated electronics were of conventional design, consisting of a high-voltage supply, a model 100 linear amplifier, and two scale-of-54 scalars operated in series.

SCINTILLATION COUNTER

Because it was hoped to take advantage of the difference in energy between the background and the radiation under investigation, a scintillation spectrometer was constructed. Although this discrimination did not prove feasible, the scintillation counter was used for the measurement of all induced isomeric transitions.

Figure 5 shows a block diagram of the counter setup.

The counter itself was an RCA type 5819 phototube with an optically clear anthracene crystal 0.5 by 2 by 4 centimeters cemented to the tube window with Canada balsam. For increased light reflection,

[illegible][illegible]

1. The number of people who are in the same line of business as the person who is being investigated.

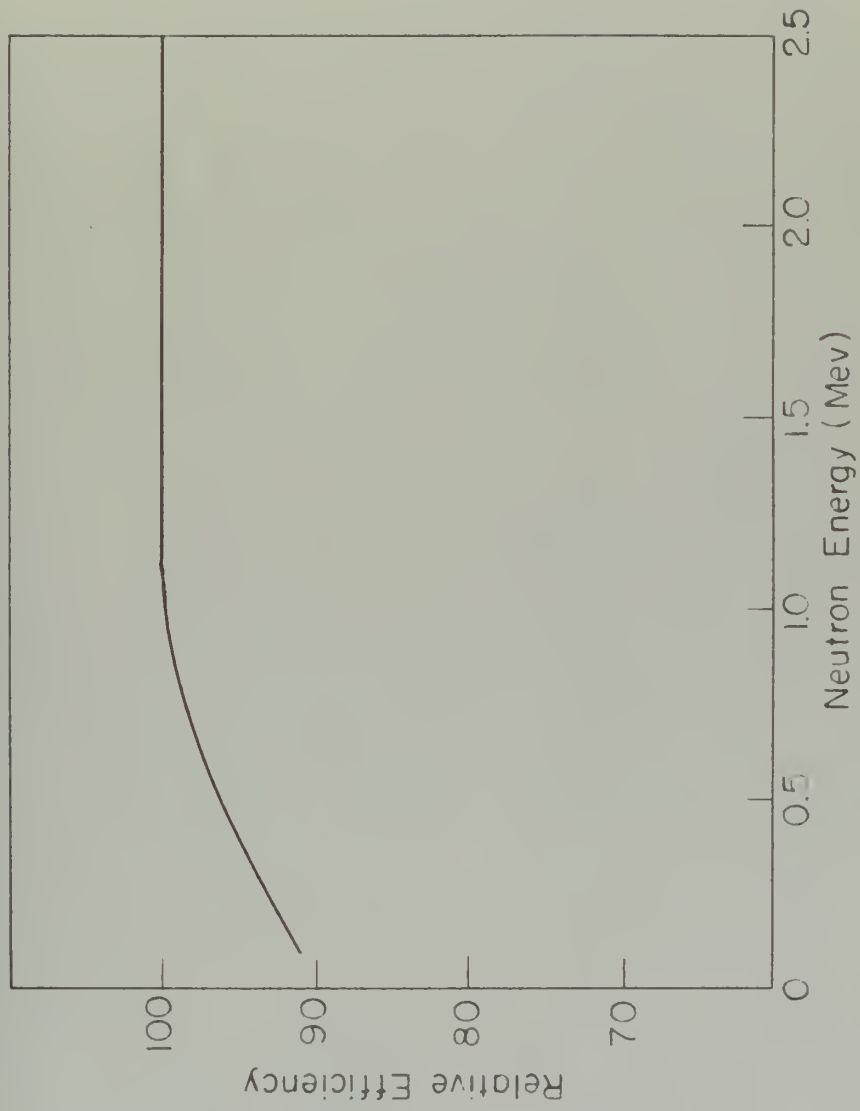


Figure 4
LONG COUNTER EFFICIENCY for DIFFERENT ENERGY
NEUTRONS

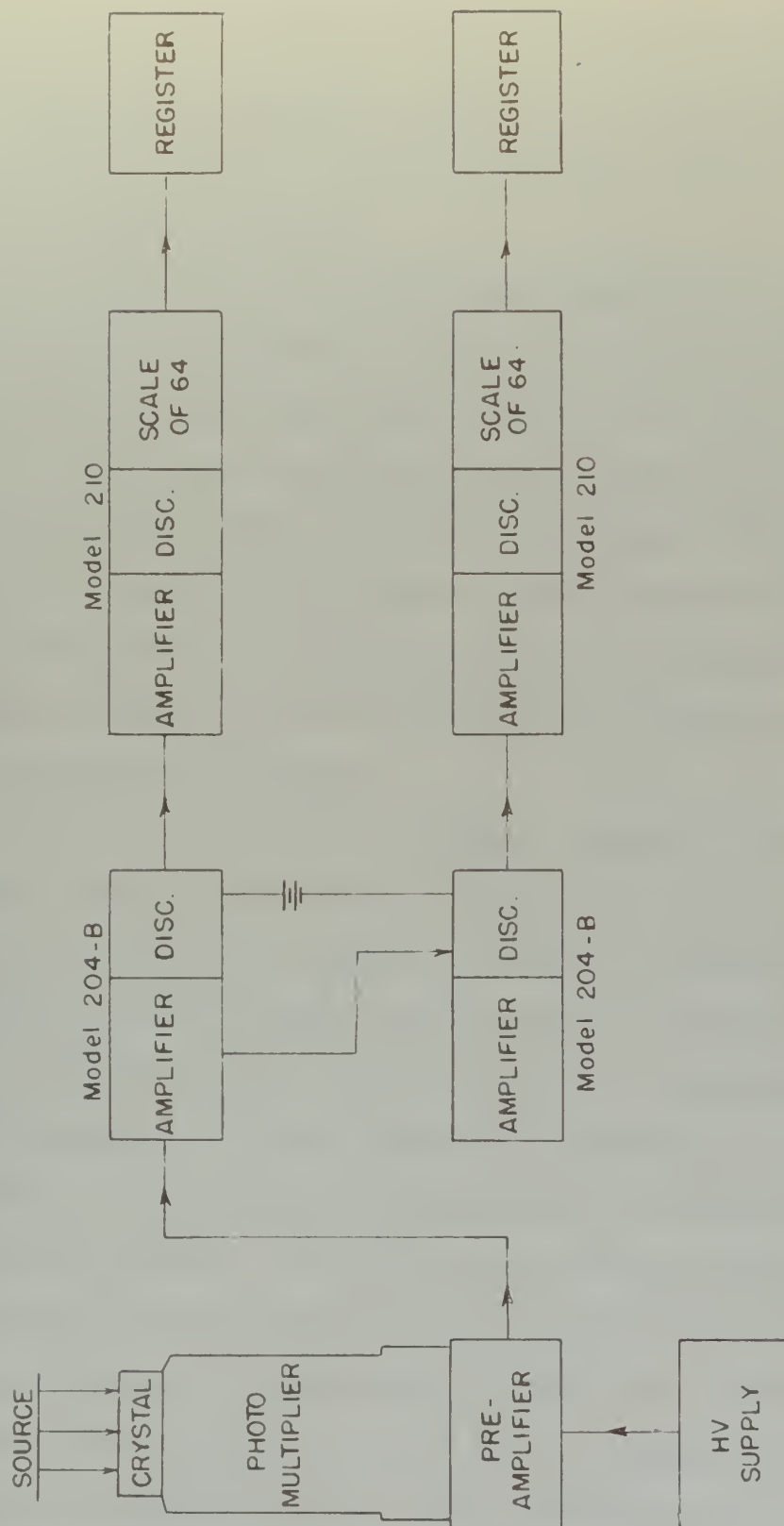


Figure 5
SCINTILLATION BETA SPECTROMETER

the crystal was covered with thin aluminum foil (1 air cm. thick). The tube and crystal were surrounded on the sides by a 0.040 inch ni metal shield for reducing stray magnetic fields and by 3 inches of lead to reduce the radiation background.

To permit rapid changing of samples, a brass plate containing a slide was fitted on the crystal end of the lead shield. The sample, in the form of a 3-cm. diameter foil, could be placed in the slide and moved to a point opposite the crystal about 3 millimeters from it.

Because of the critical dependence of phototube pulse-height output on the dynode voltages, it is important that the high voltage be kept very constant. The high-voltage supply in the Atomic Instruments Model 105 Decade Scalar was used, and the voltage was monitored with an electrostatic voltmeter.

The phototube signal was fed through a one-tube cathode follower preamplifier into an Atomic Instrument Model 204B linear amplifier. This amplifier contains an integral pulse-height discriminator. Since the discriminator bias control furnished with the amplifier did not permit reproducible settings, it was replaced by a fifteen-turn Beckman helipot. To obtain a differential discriminator, the integral discriminator from another 204B amplifier was fed in parallel with the first discriminator. The voltage for the bias of the second discriminator was taken from the first with a mercury cell in series to provide the window. This system made the window error equal to the variation in the voltage of the series battery instead of the tracking error in two potentiometers. The two discriminators were fed into separate scalars, and the differential count obtained by subtraction.

the crystal was covered with this substance (left it for 24 hours). The
 tube and crystal were surrounded on the sides by a 20 mm. layer of water.
 shield for reducing stray magnetic fields and by a layer of lead to
 reduce the scattered background.

The crystal was mounted on a goniometer, a brass plate containing a
 slide was fixed on the crystal and of the lead shield. The sample
 in the form of a 3 mm diameter disk, held in place by the slide and
 moved to a point opposite the crystal about 1 millimeter from it.

Because of the crystal dimensions of 10 mm. by 10 mm. by 10 mm.
 except on the square surface, it is important that the crystal surface
 be kept very constant. The high-voltage supply in the crystal holder
 made about 100 volts below was used, and the voltage was measured
 with an electrostatic voltmeter.

The detector signal was fed through a one-inch cable to a
 preamplifier tube on a shielded base (about 100 mm. diameter).
 This section contains an integral pulse-height discriminator. When
 the discriminator bias control (connected with the multiplier and not
 with the discriminator) was set, it was adjusted by a 100-ohm
 potentiometer. To obtain a differential discriminator, the low-
 level discriminator bias control was set to match the
 level discriminator bias control. The voltage for the bias of the second dis-
 criminator was taken from the first with a voltage divider in series to
 provide the signal. This system with the window error could be the
 variation in the voltage of the window divider instead of the level
 error in two discriminators. The two discriminators were the same
 except for the discriminator bias control, and the discriminator bias control was obtained by adjustment.

Figure 6 shows the pulse-height distribution obtained from a Cs^{137} source. The cesium activity contains a strong 630-kev conversion electron line which provided an excellent energy calibration point for the counter.

It was a most interesting and profitable trip. The main results of the trip were the following:

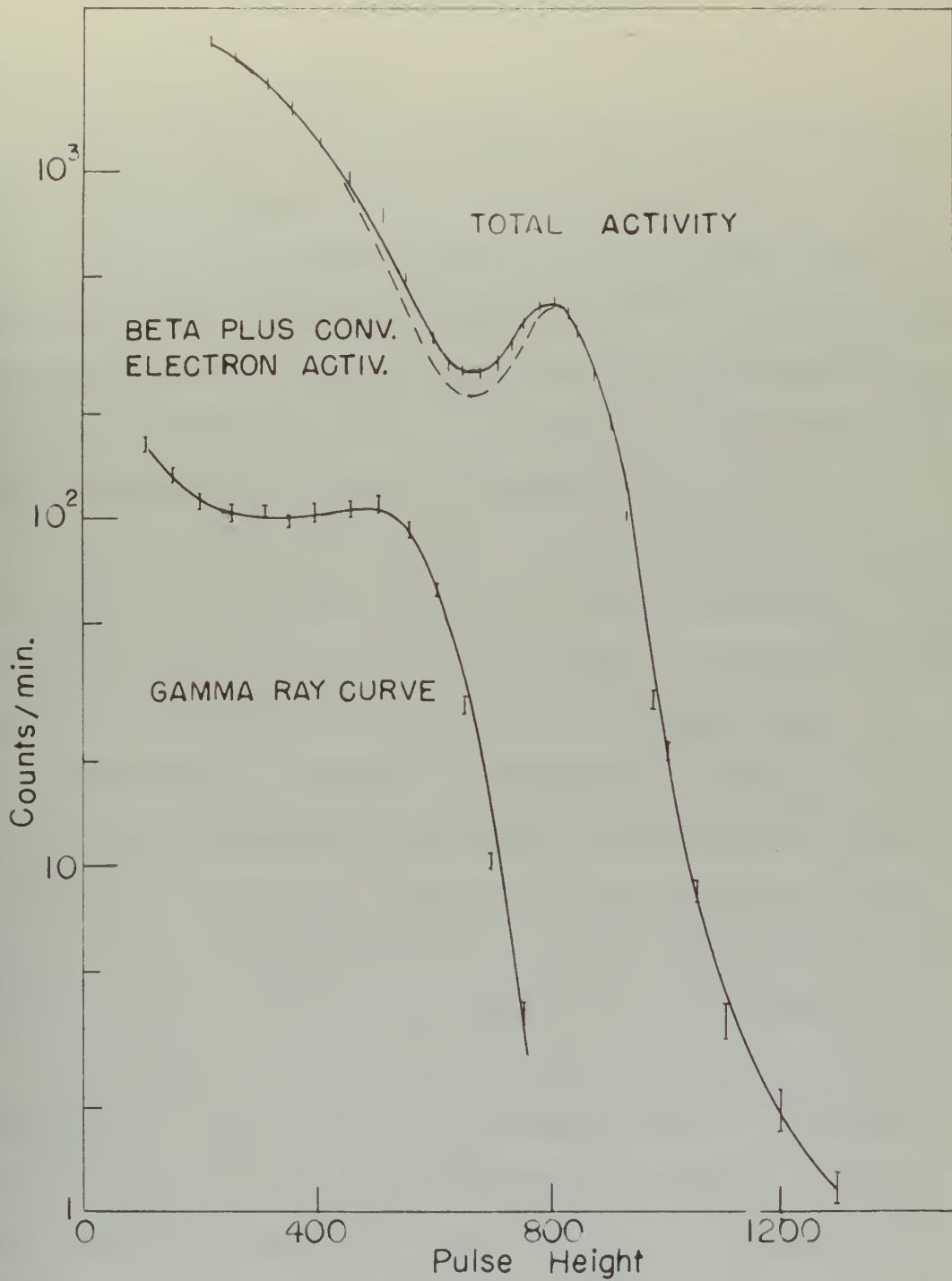


Figure 6

Cs^{137} DIFFERENTIAL PULSE HEIGHT DISTRIBUTION SHOWING 630 Kev CONV. ELECTRON LINE USED FOR ENERGY CALIBRATION.

IV. INVESTIGATION OF GOLD

HISTORY OF Au¹⁹⁷* EXCITATION

In an article surveying the field in 1941, Mattauch²³ observed that no cases of isomerism had been found for nuclei with a spin $1/2 < I < 9/2$. This observation became known as "Mattauch's second rule." The first exception to be found was Au¹⁹⁷ with a spin in the ground state of $3/2$. Later results indicate that the metastable state does not decay directly to the ground state, so that technically even this activity does not violate this empirical "rule."

Wiedenbeck²⁹ in 1945 reported a 7-second metastable state in gold excited by bombarding a gold cathode counter with x-rays. By using a thick gold target and aluminum absorbers, he was able to determine that the approximate energy of this level was 250 kev.

Subsequently, Wiedenbeck³⁰ investigated the quantum excitation curve for gold using the x-rays produced by monoenergetic electrons incident on a thick target. His results showed a series of levels starting with a threshold at 1.22 Mev and spaced about 0.4 Mev apart. The curve obtained and the levels assigned are shown in Figure 7.

In addition, he reported having excited this level by means of fast neutrons. Using a source of continuous x-rays, he obtained a broad distribution of neutron energies from the Be⁹(γ ,n) reaction. The experimental arrangement is shown in Figure 8. In the region of maximum neutron energy from 1.22 to 1.3 Mev, the counting rate in his gold counter was considerably above background and increased linearly with energy in much the same way as with x-rays. He concluded that the

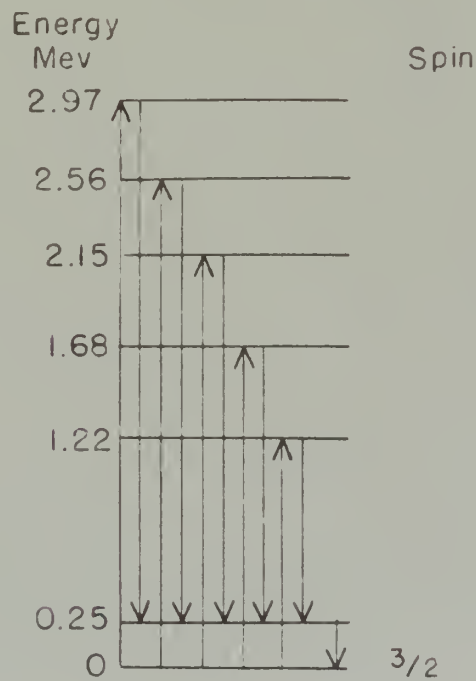
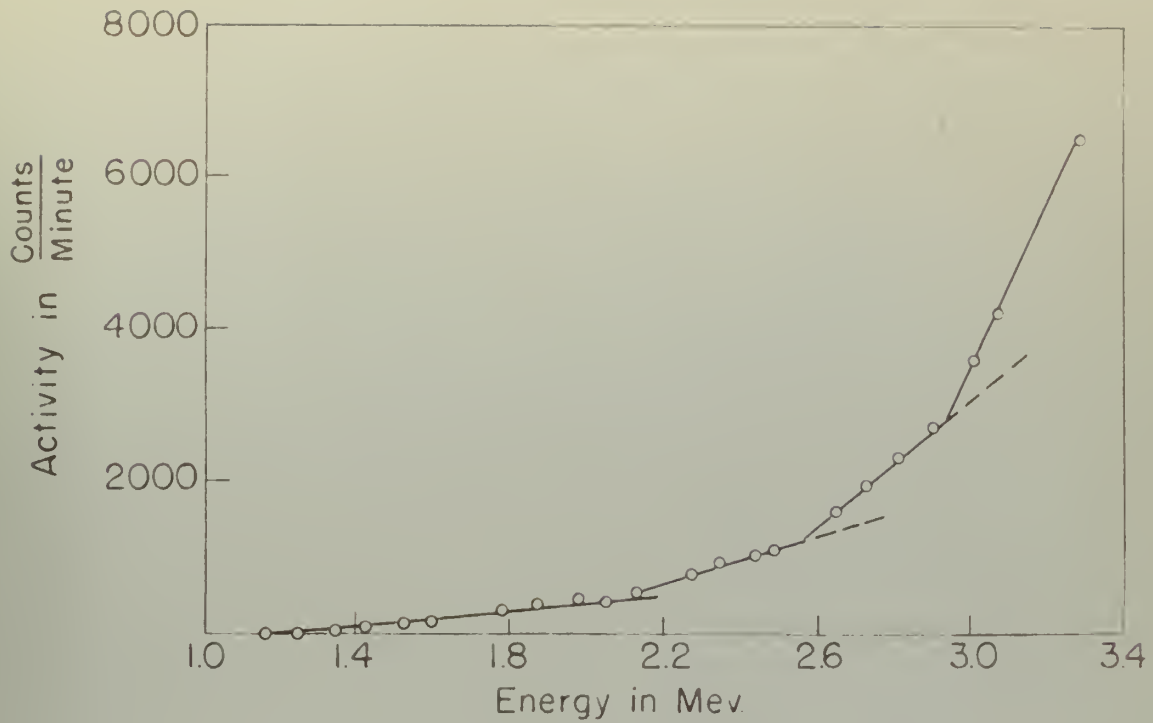


Figure 7
 QUANTUM EXCITATION CURVE for Au^{197*} and LEVELS
 OBTAINED by WIEDENBECK. (From Phys. Rev. 68, 1 (1945))

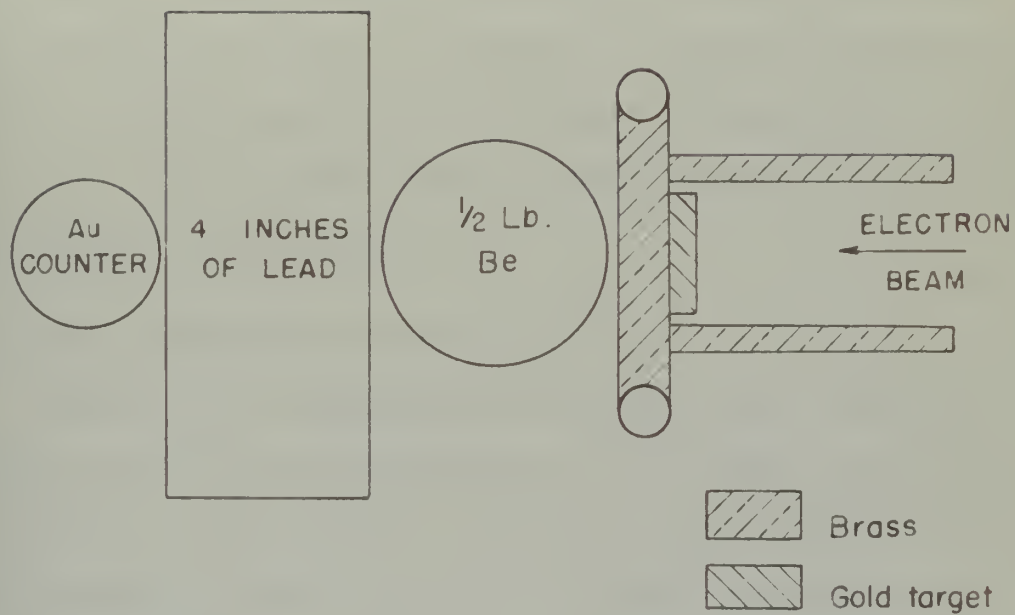


Figure 8

ARRANGEMENT USED by WIEDENBECK to INVESTIGATE NEUTRON EXCITATION of Au^{197} .

(From Phys. Rev. 68, 1 (1945))

threshold for neutron excitation was the same as for x-rays. His investigation, however, was limited to the energy interval close to the photon threshold by maximum neutron energy available.

In the determination of the excitation function, using either quanta or particles having a continuous energy distribution, the levels only appear as slight changes in the slope of a curve; hence, the energy assignment and often the existence of the level are uncertain and may even involve a subjective decision. If monoenergetic sources are used, the breaks in the curves are much sharper, and the results are considerably more certain.

DESCRIPTION OF THE ENERGY LEVELS IN Au^{197}

Because of the complexity of its excited states, the various levels in gold have been the subject of a great deal of nuclear investigation. On the other hand, gold has the great simplification of being monoisotopic with a mass of 197. Another advantage in the present investigation is the absence of any short-lived neutron capture reaction; the only one known leads to Au^{198} with a 2.7-day half-life.

The levels in Au^{197} fall into three apparently independent groups,³¹ with no evidence of transitions between any of the groups, as shown in Figure 9. The levels are shown in three groups to emphasize the improbability of transitions between the groups.

The first group consists of a level at 77 keV that is reached by K-capture from the lower 65-hour isomeric state of Hg^{197} . It has been determined that this 77-keV level does not have an appreciable lifetime.

the first group consists of a few of the most important
 the second group consists of a few of the most important
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THE FIRST GROUP

The first group consists of a few of the most important
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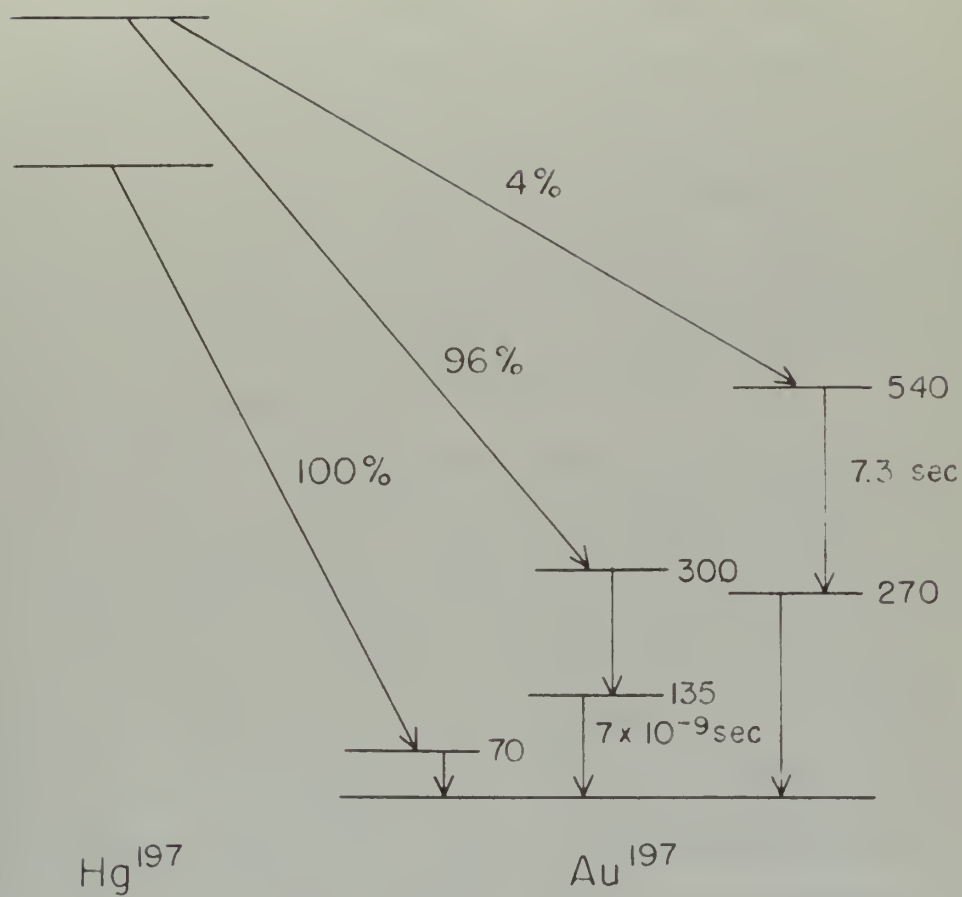


Figure 9

LEVELS in Au^{197} as DETERMINED by the
 DECAY of Hg^{197} (Energy in Kev)

The upper isomeric state in mercury with a half-life of 25 hours also decays by K-capture. In 96 percent of the cases, the decay is followed by the emission of a 165-kev gamma-ray to a 7×10^{-9} second isomeric level which in turn decays by the emission of a 135-kev gamma-ray^{32,33}. Assuming this transition to be electric quadrupole or magnetic dipole, the Weissacker hypothesis predicts a half-life for the gamma-transition of 6.5×10^{-9} seconds. This neglects the effect of internal conversion on the lifetime, but the calculations of Rose et al.³⁴ indicate that the conversion coefficient will be less than unity and, hence, not sufficient to reduce the predicted half-life for a higher-order transition below a few microseconds.

The remaining 4 percent of the nuclei in the 25-hour mercury decays to a 7.3-second level which in turn decays by the emission of two internally converted gammas both of about 270 kev³⁵. Originally, the lower level was reported to be only of the order of 70 kev³⁶, and this value is found in much of the literature. This determination was made by means of an absorption of the conversion electrons in aluminum by Frauenfelder, Huber et al. However, the reinvestigation by members of this group, using a magnetic spectrometer, revealed a coincidence rate in the K-conversion line for 270-kev gammas that was about twenty times the chance coincidence rate. The ratio of K- to L-conversion was also determined as 3.4.

Using the Weissacker calculation of the half-life with an energy of 270 kev, the following values are obtained for several different changes in nuclear angular momenta:

The report contains a table showing the results of the tests conducted on the various specimens of the material. The table is as follows:

Specimen	Test	Result
1	Tensile	100,000 lbs.
2	Compression	150,000 lbs.
3	Flexure	120,000 lbs.
4	Impact	110,000 lbs.
5	Shear	130,000 lbs.

The results of the tests show that the material is capable of withstanding a tensile stress of 100,000 lbs., a compressive stress of 150,000 lbs., a flexure stress of 120,000 lbs., an impact stress of 110,000 lbs., and a shear stress of 130,000 lbs. These results are in good agreement with the values given in the specification.

TABLE II

Δ	Parity Change	Half-Life (Seconds)	Electric pole	Magnetic pole
2	no	2.6×10^{-10}	2^2	2^1
3	yes	5.4×10^{-5}	2^3	2^2
4	no	22.	2^4	2^3
5	yes	1.6×10^7	2^5	2^4

The symbol Δ is introduced to characterize an electric 2^Δ or magnetic $2^{(\Delta-1)}$ transition, both having the same probability under the Weissacker hypothesis.

From this table, the transition from the 7-second state is evidently of the $\Delta = 4$ type and, hence, accompanied by no change in parity. However, it is not possible to determine from the half-life whether the transition is an electric 2^4 pole which involves a spin change of 4 or is a magnetic 2^3 pole with a spin change of 3. The ambiguity is removed by the conversion coefficient data which show the transition to be of the magnetic type³⁷.

Frauenfelder³⁸ reports that the state to which this isomeric level decays has no appreciable lifetime, and Segre and Helmholtz³⁹ set this lifetime at less than 1 microsecond. Later in this section, it will be shown that this transition cannot be of the $\Delta = 2$ type, because of parity considerations. In view of the crudeness of the theoretical calculations of lifetimes, the possibility cannot be ruled out that this is a $\Delta = 3$ transition with a predicted lifetime of 54 micro-

seconds. The other possibility corresponds to an electric dipole transition which in general does not occur because of the smallness of the electric dipole moment.

Figure 10 shows the energy levels involving the 7-second metastable level, together with the parity and spin assignments that may be made from neutron excitation of the isomer.

Because the decay of the 7-second level directly to the ground state is not observed, although the energy is twice that for decay in cascade, this transition must have a Δ value greater than 4. However, this level may be excited directly by neutrons, as is shown later in this chapter.

A comparison with theory of the experimental cross section for the production of the isomeric state shows that the neutron cannot carry in more orbital angular momentum than $\ell = 3$. The shape of the excitation curve shows that the neutron is emitted with $\ell = 0$. To be added to this value of $\ell = 3$ is the contribution of the neutron spin which may be -1, 0, or 1, depending upon the orientation of the spins of the incident and emitted neutrons relative to the orbital momentum. Thus, a $\ell = 3$ neutron can excite a level with a spin differing from the ground state by 2, 3, or 4 units. This reaction gives a change in parity of the excited state due to the ℓ value being odd for the incident neutron and even for the emitted neutron.

The spin changes of 2 or 3 between the ground and the excited state with a parity change both give a lifetime corresponding to $\Delta = 3$ which is much too short. The only other possibility, the spin change 4

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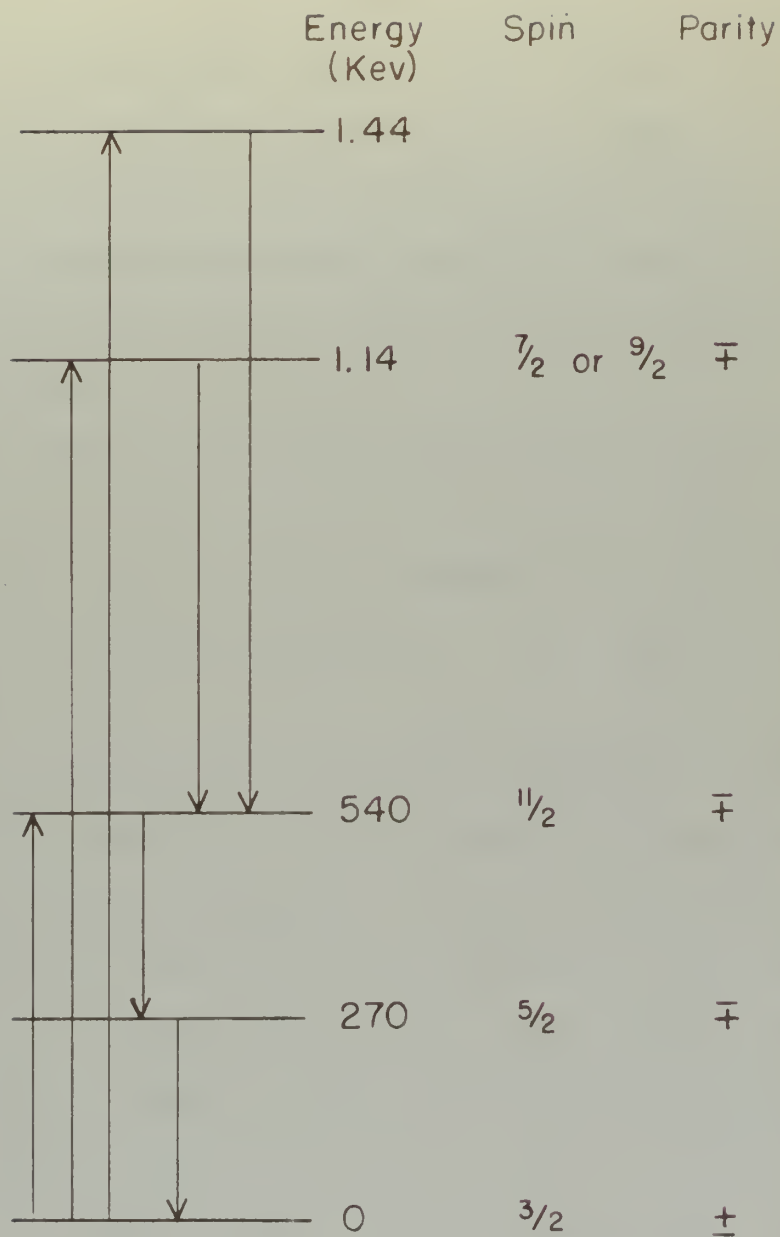


Figure 10

LEVELS EXCITED by NEUTRONS in Au^{197} which
DECAY to the METASTABLE STATE at 540 Kev.

with a parity change, however, gives a magnetic 2^4 -pole transition with a lifetime approximated by letting $\Lambda = 5$. Thus, the spin and parity relationship between the ground state and the 7-second metastable state are uniquely determined, and it is shown that the neutron enters the Au^{197} nucleus with orbital angular momentum of 3 and its spin parallel to the orbital axis, and it is emitted with angular momentum zero and its spin reversed.

Since the isomeric transition involves no parity change, the level to which it decays must also have a parity different from the ground state. This rules out the possibility of a transition from this state to the ground state having Λ even, and in particular it excludes $\Lambda = 2$.

Two other levels at 1.14 and 1.14 Mev are found that may be excited by fast neutrons of < 2 -Mev energy. However, at these energies, the incident neutron may carry in even more angular momentum. At the same time, there is no lifetime requirement on the state, except that transitions to the metastable state are quite probable. Furthermore, the shape of the excitation curve for each level individually cannot be determined because of the uncertainty in the contribution of lower levels to the isomeric state. For these reasons, little can be deduced from the neutron excitation curve about these states other than their energy.

It is of interest to compare the energies determined by neutron excitation with those Wiedenbeck has reported for quantum excitation. In Chapter II it has been shown that any level excited by a quantum should be excited to an appreciable degree by neutrons. The level

excited by 1.14-Mev neutrons apparently corresponds to the threshold found for 1.22-Mev quanta. The agreement is considered satisfactory because of the difficulty of deducing this threshold accurately from the small slope between the quantum excitation curve and the axis. However, no evidence is found with neutrons of the level reported for quanta at 1.68 Mev.

The fact that the 1.14-Mev level can be excited by quanta enables a determination to be made of its spin relative to the ground state. The spin change of $\frac{1}{2}$ from the ground state to the metastable state with a parity change must be accomplished by the absorption of a quantum plus the emission of another.

There are two possibilities with high enough probabilities for the process to occur. The excitation is either by an $\ell = 2$ yes (magnetic quadrupole) radiation followed by decay to the isomeric state by $\ell = 2$ no (electric quadrupole) radiation or excitation by $\ell = 3$ yes (electric octupole) radiation followed by $\ell = 1$ no (magnetic dipole) radiation. It is impossible to separate these since they both have the same probability according to the Weissacker hypothesis. However, the decision as to whether a transition corresponds to the excitation of the higher level or the decay to the metastable level is unambiguous because the decay to the metastable level must be able to compete successfully with the decay to the ground state in spite of the higher energy involved in the latter.

Unfortunately, the Wiedenbeck data do not permit even an inference regarding the probability of exciting the neutron level at 1.14 Mev with quanta.

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DESCRIPTION OF EXPERIMENTAL SETUP

Measurement of the neutron excitation function of Au^{197*} was complicated by two factors:

1. The cross section for the excitation process is so small that it was necessary to sacrifice part of the energy resolution to obtain activities large compared to the background; and,
2. The half-life of the isomeric state is only 7 seconds.

Figure 11 shows the experimental setup used for the gold investigation.

The technique used involved irradiating the sample in the form of a foil with a known flux of monoenergetic neutrons for 15 seconds, cutting off the neutron beam, transferring the foil to a scintillation counter, and counting the activity for 15 seconds.

The gold foil was 3 cm. in diameter and 0.1 mm. thick with a mass of 1.420 grams. To permit rapid transfer of the foil from the irradiating position to the counting position, the foil was mounted in a depression at one end of a bakelite slide 1/8 inch thick by 1-1/2 inches wide by 3 feet in length. This slide was carried in a slot in the shield of the scintillation counter, which has been described previously. When the slide was against the stop at one end, the foil was accurately positioned next to the crystal of the counter. When it was against the stop at the other end, the slide was held by aluminum fingers in front of the beam to eliminate errors due to the flexibility of the slide. This arrangement was necessary to keep the counter out of the region of high neutron flux which would lead to high background counting rates for

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Source: *Journal of the American Statistical Association*, 1977, 72, 1, 1-11.

• *continued*

The authors are indebted to J. H. Duerksen for his assistance in the preparation of this manuscript.

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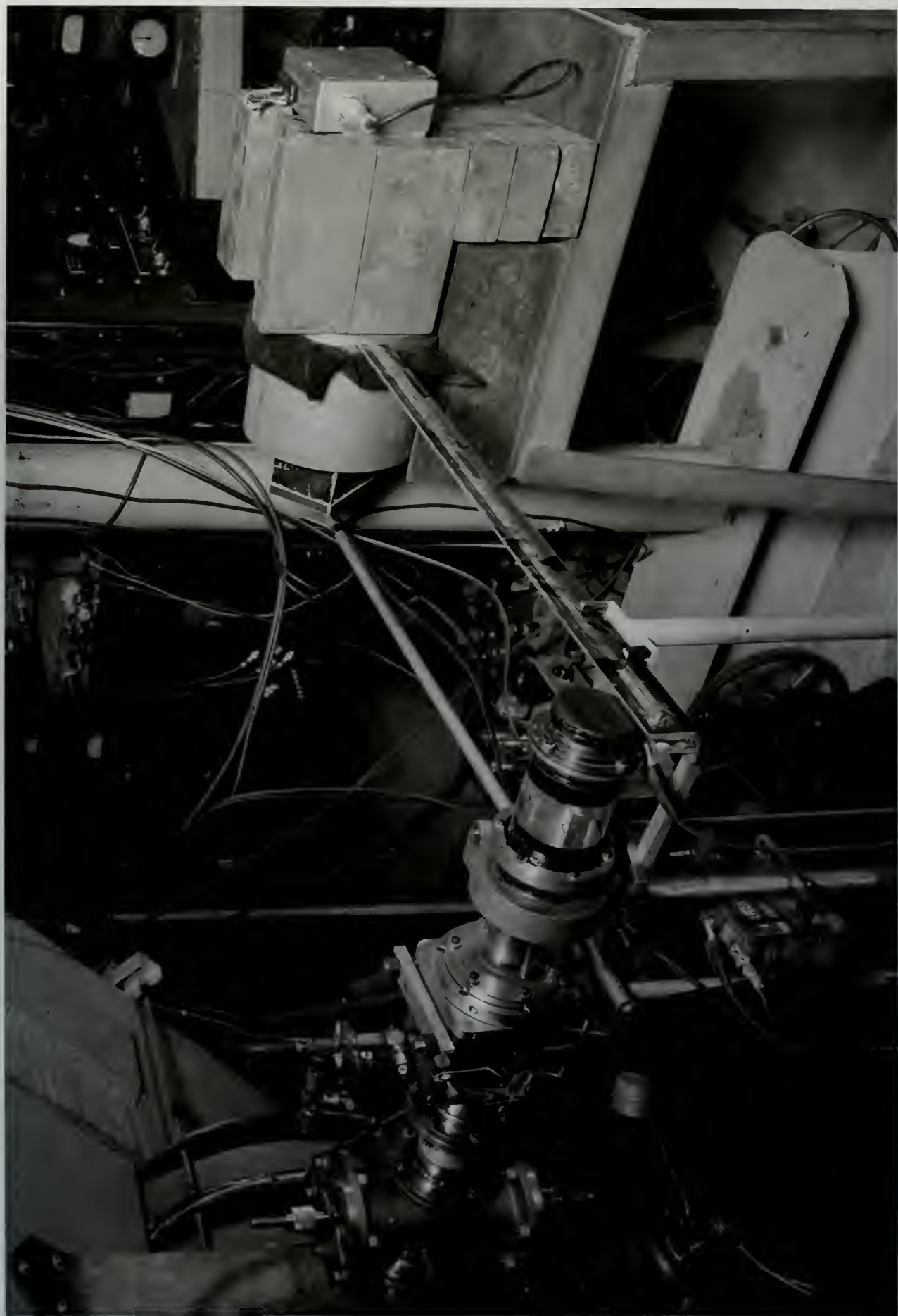
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It is possible that this study will lead to high frequency monitoring rather than

Figure 11. Arrangement for irradiating the gold foils, showing the bakelite slide and the two counters with their lead shields.



a short time after the beam was cut off. It also permitted shielding of the counter on all sides from the stray gamma-ray background that was present even when the beam was not striking the target.

During the counting periods, the neutron beam was cut off by closing a tantalum-backed flap valve in the proton beam. This valve, unfortunately, was located near the target instead of at the entrance slit of the deflection chamber. For this reason, slight wanderings resulted in the beam striking the chamber walls which caused large fluctuations in the background during the counting period.

To measure the background while the counting was taking place, another scintillation counter with characteristics as nearly as possible like the counter for the gold was set up. This new counter was placed beside the other with the same amount of shielding but also was shielded from the gold activity. The scalars for both counters were arranged so that they were controlled by the same switch. Test activation runs were then made which were identical to the true runs, except for the omission of the gold from the slide. These runs showed the backgrounds of the two counters under these conditions to be the same within a statistical uncertainty of about 5 percent. Since the background over most of the curve was of the order of 10 percent or less of the induced activity, the matching was considered satisfactory.

The neutron flux for these runs was determined by the long counter previously described. The counter was set up one meter from the target at zero degrees. The number of counts during each irradiation was recorded and corrected for the effect of neutron energy on the counter efficiency.

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The exposure of the foil was initiated by the operator's pulling a cord with his left hand which opened the flap valve and allowed the proton beam to strike the target. After 15 seconds, the cord was released and the flap was closed by a counterweight. At the same time, the operator pulled a second cord with his right hand which moved the slide into position, and with his left hand, he started the counter. This arrangement prevented the accidental starting of the counter before the beam was off. With practice a rhythm was established that permitted starting the counting period one-half second after the end of the irradiation period with a variation of less than one-fourth of a second. This variation corresponds to a random error of less than 3 percent.

TREATMENT OF THE DATA

Each observation yielded on the order of 150 to 400 counts in the 15-second counting interval with a background of about 20 counts. At each energy, five or more identical runs were made and averaged to reduce the statistical and other random errors, such as variations in the length of time between the end of the irradiation and the beginning of the counting. An analysis of the individual readings taken at the same energy shows that they agree among themselves within the statistical error. The nonexistence of appreciable errors due to slow variations was shown by the reproducibility of the data over a period of several days.

Two tests were made to check that the 7-second isomeric state of gold was being counted. A decay curve of the radiation was taken by

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noting the number of counts at the end of each 5 seconds. This method was subject to large statistical and human errors, but the results were consistent with a half-life between 5 and 10 seconds. Since there are no other known activities which could be excited in gold with this half-life and since the activity was not found when the gold was removed from the slide, this was considered sufficient to identify the activity.

In addition, a pulse-height distribution curve was run to identify the energy. This curve also suffered from poor statistics but was consistent with the known characteristics of the isomeric state.

In the determination of the excitation curve, the energy resolution was affected by three major factors, the fluctuations in the energy of the proton beam, the thickness of the lithium target, and the dependence of the neutron energy on the angle of emission. A fourth factor, scattered neutrons, may be neglected because of their low intensity relative to the direct beam in the region in which the foil was exposed.

To increase the beam current striking the target, it was necessary to operate with the beam defining slits opened to 5 millimeters. This leads to a resolution of approximately 15 kev in the proton beam and, hence, in the neutron beam. Since the beam striking the slits was large compared to the opening, all energies in this region are approximately equally represented.

The effect of using a target of finite thickness is to give also a square number-energy distribution for the neutrons similar to that

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It is a well-known fact that the most important factor in the development of a nation is the quality of its leadership. The history of the world is full of examples of nations that have risen to greatness and others that have fallen into decay. The difference between the two is often found in the quality of the leadership that has guided them. A nation that is led by a wise and just leader will prosper and its people will be happy. A nation that is led by a foolish and unjust leader will decline and its people will be miserable. Therefore, it is essential that a nation choose its leaders carefully and wisely. The people should have a say in the choice of their leaders and should hold them accountable for their actions. Only in this way can a nation hope to achieve greatness and lasting peace.

to increase the time spent studying the subject, it was necessary to compare with the time spent on other subjects. This time is a measure of the amount of time spent on the subject and, hence, to the subject being. From the time spent on the subject compared to the amount, all needed in the subject was undoubtedly equally important.

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caused by the slit opening. The effects are therefore additive and, according to measurements on the rise of the neutron yield curve at threshold, amount to a resolution of 30 kev in this investigation.

The third factor in the resolution, the angle at the target intercepted by the foil, must be considered because of the dependence of the neutron energy on angle. For those experiments in which the cross section for the reaction is quite small, it was found necessary to sacrifice angular and, hence, energy resolution in order to place the foils close enough to get sufficient activity. Accordingly, the foils were exposed 5 centimeters from the target. They intercepted a half angle of 17 degrees for the portion of the curve above 1.0 Mev. Below 1.4 Mev the exposure was made at 4 centimeters with a half angle of 21 degrees. This corresponds to an energy spread of 25 kev at 600 kev in the forward direction and to 35 kev at 2.0 Mev.

Figure 12 shows the correction that must be applied to the maximum neutron energy to obtain the mean energy of neutrons incident on the foil. This curve was obtained by the following considerations.

For the purpose of calculating the mean energy incident on the foil, the beam may be considered uniform in intensity over the angle involved since the energies involved are well above the $\text{Li}(p,n)$ threshold.

Let the energy spread due to the slit width plus the target thickness be represented by ΔE_p and that due to the finite solid angle intercepted be ΔE_θ .

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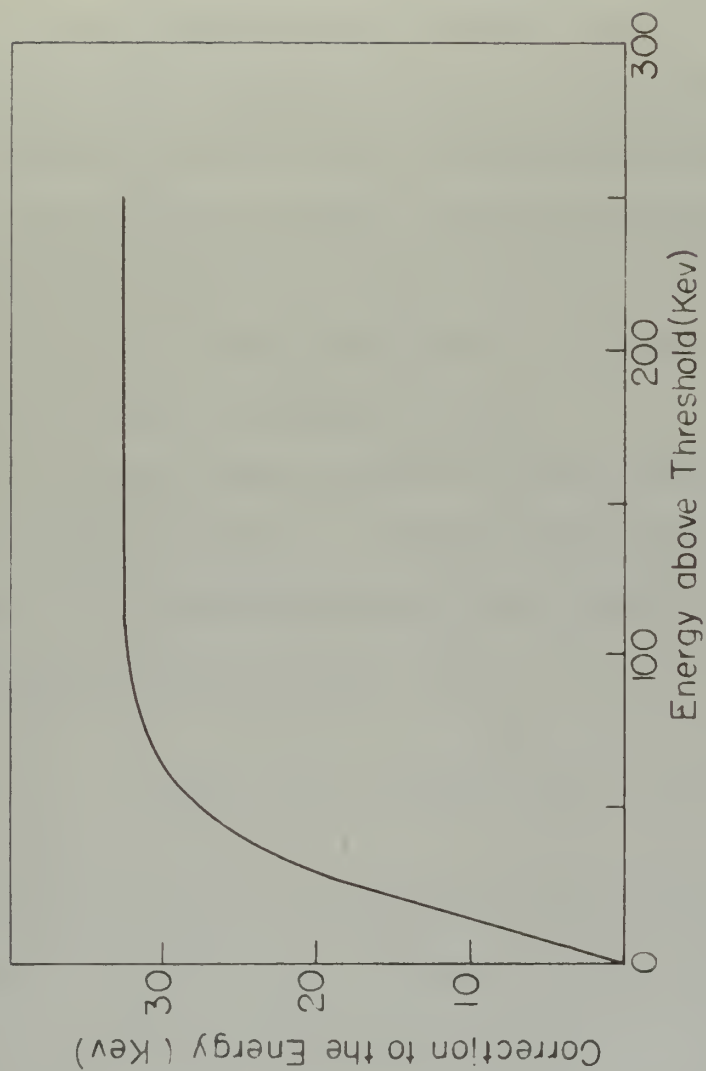


Figure 12

ENERGY CORRECTION to be SUBTRACTED from the MAXIMUM ENERGY INCIDENT on the FOIL in the VICINITY of the EXCITATION THRESHOLD.

The variation of the cross section for the production of the isomeric state may be considered constant over the energy spread ΔE_0 except at threshold. Assume for the moment that the neutron beam is monoenergetic at 0 degrees; that is, that ΔE_p is zero. Then, the activity due to each increment of energy is proportional to the area of the target subtended by the corresponding increment of angle at the source. This obviously increases linearly with energy up to the angle intercepted by the whole target foil. The mean energy in terms of the activity it produces is then calculated to be 0.707 of the total spread, ΔE_0 , below the 0 degree energy.

This result may now be integrated over the spread of neutron energies present at 0 degrees. Again assuming the cross section to be constant over the energies involved, the correction to be subtracted from the maximum energy is found to be:

$$0.5 \Delta E_p + 0.707 \Delta E_0$$

The calculation of course does not hold at threshold because of the assumptions on the effect of energy on cross section. When the maximum energy is just equal to the threshold for excitation, the correction must be zero because that is the only energy causing excitation.

Another point may be calculated by considering the case in which the lowest energy present is just the threshold energy. Again, with somewhat less justification, assume the cross section uniform over each ΔE_0 , but this time assume that the cross section over ΔE_p increases

as $(E - E_{th})^{1/2}$. A calculation similar to that preceding gives a correction equal to:

$$0.37 \Delta E_p + 0.707 \Delta E_n$$

This correction is somewhat too large due to the assumption on ΔE_n but gives a reasonable approximation of the true value.

The activity of the gold above background was divided by the neutron flux and is plotted in Figure 13 as a function of the mean energy of the neutron beam as corrected according to the preceding considerations. The upper portion of the curve taken with the foil 5 centimeters from the target was normalized to the activity expected at 4 centimeters to join with the lower portion of the curve. The two sections overlap in the energy interval from 1.0 to 1.4 Mev.

The neutron excitation curve shows the characteristic shape for three levels at 530 ± 20 kev; 1.14 ± 0.03 Mev; and 1.44 ± 0.03 Mev with a possible fourth level at about 1.77 Mev. The errors assigned are estimates of the uncertainty in the determination of the energy of the neutrons and the location of the level on the curve.

The section of the curve in the vicinity of the 1.44-Mev level was checked with higher resolution to be sure that the level was actually present. It was found that the break in the curve was readily reproducible.

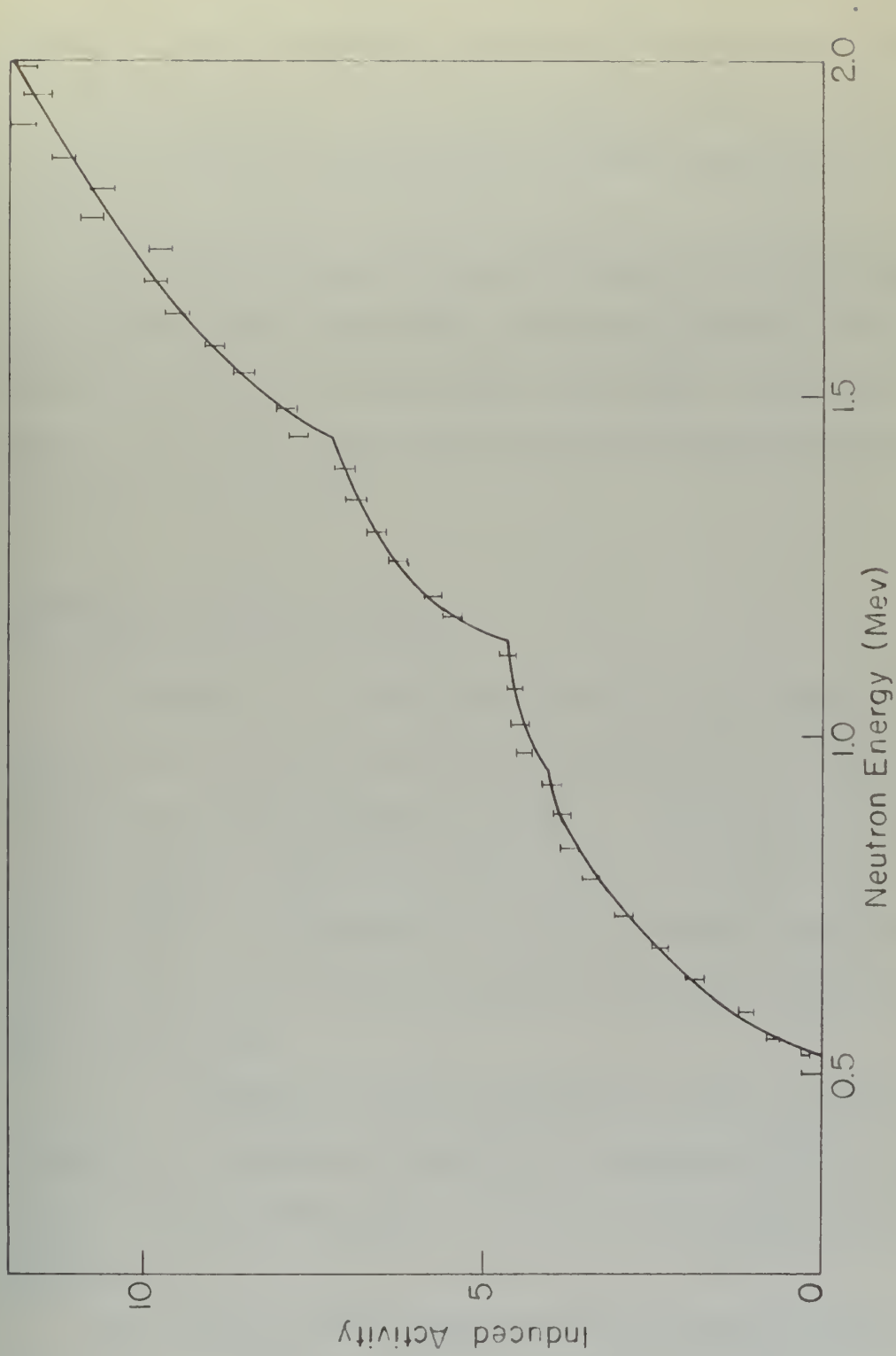


Figure 13
NEUTRON EXCITATION CURVE for Au¹⁹⁷*

The curve shows about a 10 percent rise 430 kev above threshold. Although this may be explained by statistical uncertainty, it seems more probable that this is due to the second energy group reported to be 384 kev below the main group with a relative intensity of 10 percent.

The threshold value is of special interest because it corresponds quite closely with Huber's measurement which shows the metastable level to be about 540 kev above the ground state. As a result, there appears to be little doubt that the neutrons are able to excite the metastable state directly.

COMPARISON OF THE THEORETICAL AND EXPERIMENTAL CROSS SECTIONS

A knowledge of the cross section for the formation of the metastable state is of great importance in assigning an ℓ value to the incident neutron that excited that state. A comparison of this cross section with the $\sigma_c^{(\ell)}$ cross section for the formation of the compound nucleus Au^{198} as a function of ℓ for the incoming neutron is sufficient in this case to place a maximum on the angular momentum the neutron can carry in.

The calculation of $\sigma_c^{(\ell)}$ following the method of Blatt and Weisskopf, as described in Chapter II, is carried out in the energy region of 600 to 1200 kev for $0 \leq \ell \leq 4$. The results are plotted in Figure 14.

Determination of the experimental cross section involves knowledge of the neutron flux on the gold foil and the efficiency of the

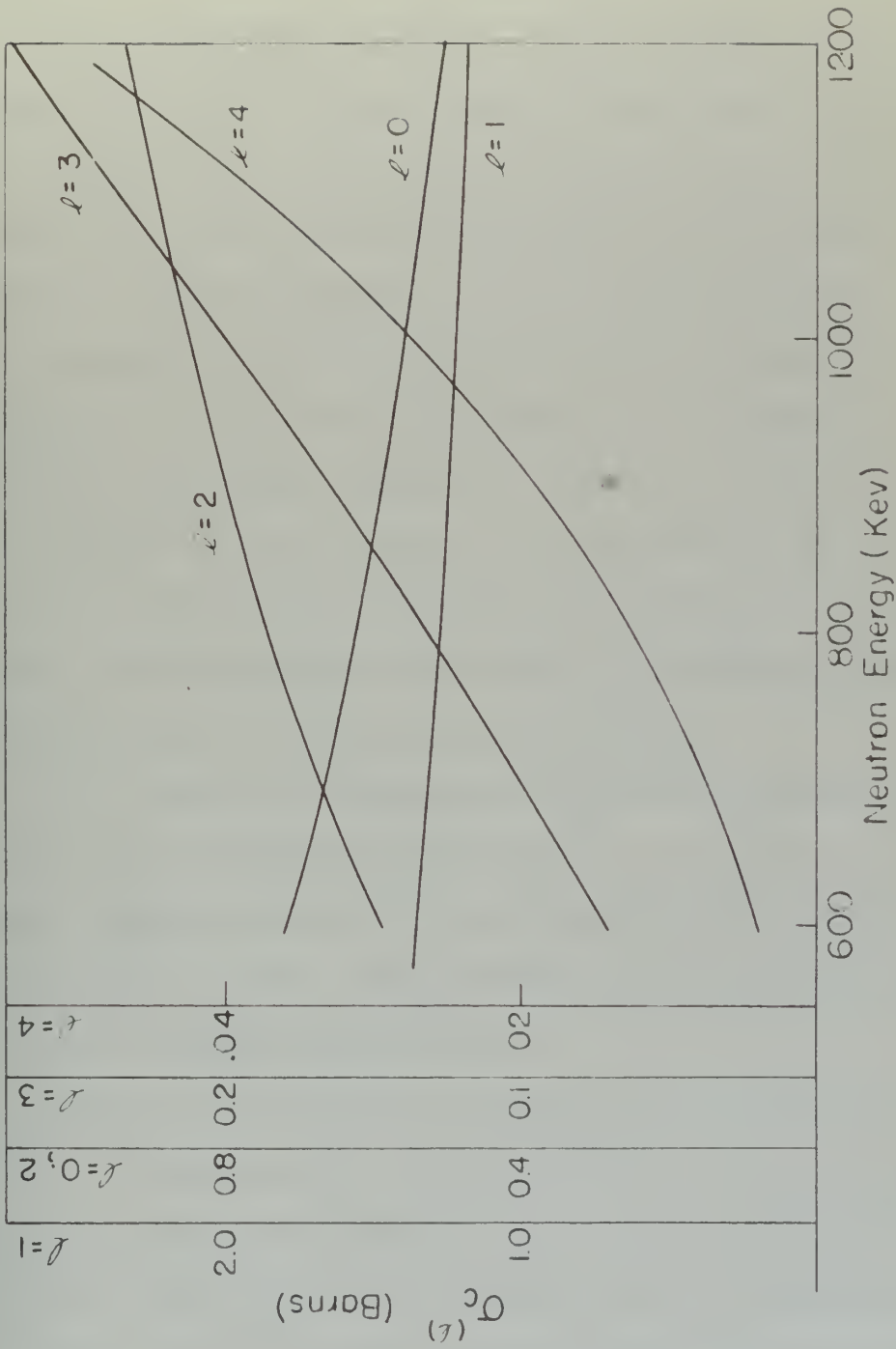


Figure 14

CROSS SECTION for FORMATION of the COMPOUND
NUCLEUS $\text{Au}^{197} + n \rightarrow \text{Au}^{198}$
(Note the different scales for different l 's)

counter for counting the excited state. The largest error in this cross section will result from the difficulty in evaluating the second factor.

In the data plotted on Figure 13, the ordinate is the number of counts of the Au^{197*} activity divided by the number of register counts from the long counter when operating through a scale of 1024 scalar. Willard⁴⁰ has determined that the flat part of the neutron yield curve this counter at 0 degrees and one meter from the target gives 60 counts per microcoulomb of protons per kev of target thickness, the target thickness being measured by the initial rise of the neutron yield curve with good proton energy resolution. From the data of Hanson, Taschek, and Williams⁴¹ on neutron yield, a 40-kev target in this energy region gives 4.35×10^6 neutrons per microcoulomb per unit solid angle at 0 degrees. Therefore, each count on the long counter corresponds to 1610 neutrons per unit solid angle at 0 degrees.

When exposed at 4 centimeters, the 3-centimeter foil intercepts a solid angle of 0.123 steradians. Hence, each count on the long counter indicates that the foil has been traversed by 232 neutrons; or, taking into account the scaling factor, each register count stands for 2.38×10^5 neutrons through the foil.

The area of the foil is 7.07 square centimeters, and the mass is 1.420 grams corresponding to 4.34×10^{21} Au^{197} nuclei. Therefore, an ordinate of unity on the plot in Figure 13 is equivalent to a cross section of 0.0069 barns if it is assumed that the scintillation counter counts every gold nucleus excited. This, then, is the minimal value

These people will remain free and friendly to the world
because for example the world state. The largest group in this

[illegible]

There is a lot of talk about the "new" world, but it is not new at all. It is the same old world, only with a different name. The world has always been the same, and it always will be. The only thing that has changed is the way we look at it. We have a new name for it, but it is still the same old world.

[illegible]

that the cross section can have. On Figure 15, the minimal cross section for the production of the metastable state in the vicinity of the threshold is plotted as a function of energy, together with the cross section for the formation of the Au^{198} compound nucleus by an $\ell = 4$ neutron, $\sigma_c^{(4)}$. It is seen that $\sigma_c^{(4)}$ is even less than this minimum over part of the region. This reaction can therefore be definitely ruled out as the method of exciting the 7-second state directly.

Of course, the minimal value computed is too low by a factor equal to the reciprocal of the overall efficiency of the counter for counting the metastable state. Since the foil thickness is over 10 times the range of the conversion electrons, their contribution to the efficiency may be neglected, and the counting rate may be considered to come entirely from the gold x-rays produced on their emission and from the unconverted gamma-rays. Due to the resolution of the scalar following the counter, the decay of the nonisomeric 270-kev state cannot be separated from the isomeric state.

Using the mass-absorption coefficient for the gold K-series x-rays of $3.40 \text{ cm}^2/\text{gm}$, it can be calculated that in the one-dimensional case of a source 0.1 millimeters thick, the self-absorption is 0.73. The attenuation of this radiation in the crystal may be calculated to be 0.76, corresponding to an absorption of 0.24.

For the 270-kev unconverted radiation, the self-absorption is small and the crystal absorption 0.20. This gives an efficiency for counting the x-rays from this source of about 0.13 and for the gamma-rays about the same figure. So far, it has been neglected that only

The first of these is the fact that the
 H_2O content of the soil is not
 constant, but varies with the
 amount of water available to the
 plants. This is due to the fact
 that the plants take up water
 from the soil, and the amount
 of water taken up depends on
 the amount of water available
 to the plants. The amount of
 water available to the plants
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 in the soil, and the amount of
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 circular argument, and it is
 not possible to determine the
 amount of water available to
 the plants without knowing the
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 The amount of water in the
 soil depends on the amount of
 water available to the plants,
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 the plants without knowing the
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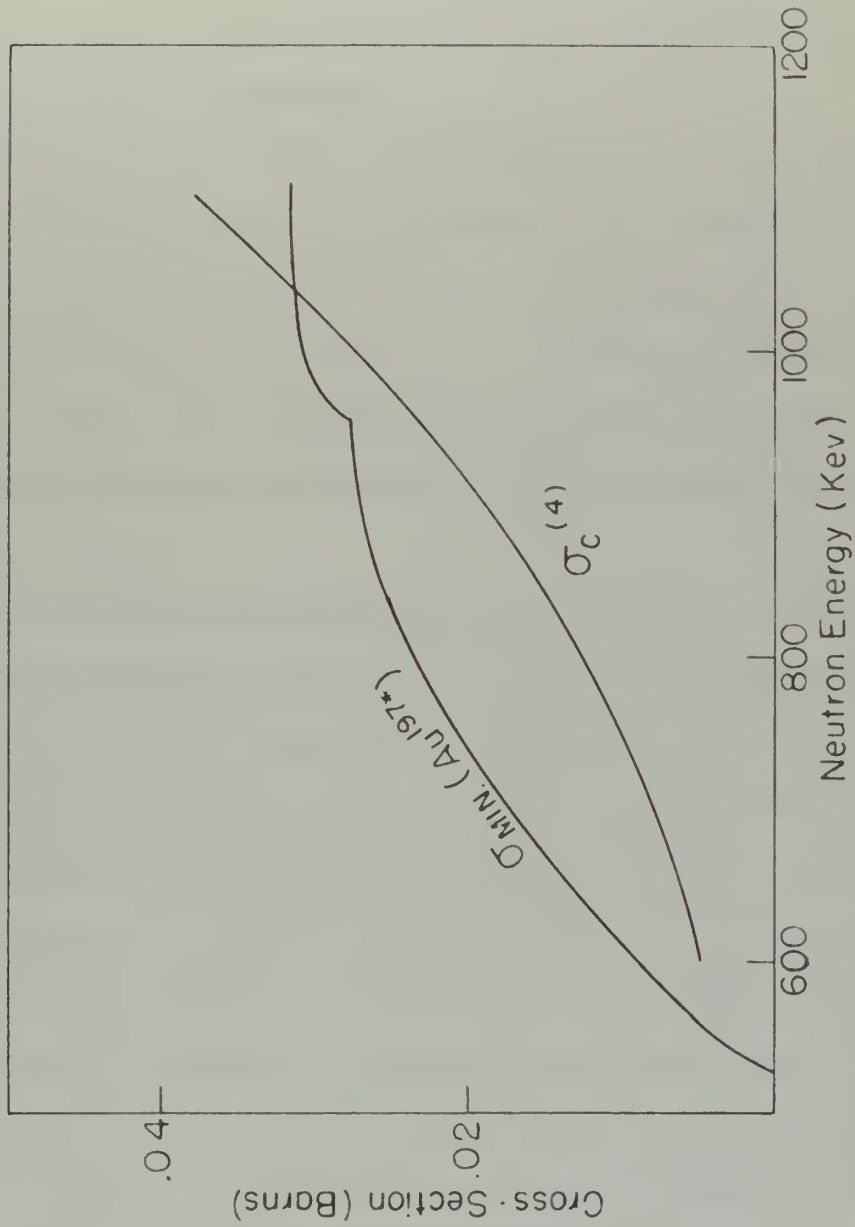


Figure 15
COMPARISON of MINIMAL CROSS-SECTION for PRODUCTION of Au^{197*} with $\sigma_c^{(4)}$ SHOWING THAT $\sigma_c^{(4)} < \sigma_{\text{min}}$ for the INCIDENT NEUTRON.

half of the radiation is directed toward the counter and that two decays take place for each metastable nucleus. However, these factors will cancel each other to the extent of the uncertainty involved in this calculation. Softer x-rays, notably the L-series, have also been neglected in this approximation.

There is another factor to be considered, the efficiency of the photomultiplier for counting scintillations in the crystal. This has been checked by comparing the counting rate in this counter from a thin C^{14} source with that obtained in a 50 percent geometry windowless gas-flow counter. It is found that this factor, which also corrects for the scintillation counter geometry being slightly less than 50 percent, is 0.90.

From all these considerations, it is estimated that the scintillation counter counts about 16 percent of the Au^{197*} nuclei decaying during the counting interval. The bias on the discriminator, however, was set so that 10 percent of the counts were not recorded, and the counting period of two half-lives was such that only 75 percent of the nuclei decayed during the period. Therefore, the correction to the cross section should be increased by a factor 1.48, giving a final value of 0.24. Because of the necessary crudeness of this calculation, it is considered that it may be in error by as much as 50 percent.

In Figure 16, the cross section for the formation of the metastable state in the vicinity of the threshold is plotted with this factor applied to the minimal cross section. A plot of the cross section for the formation of the compound nucleus by an $\ell = 3$ neutron is

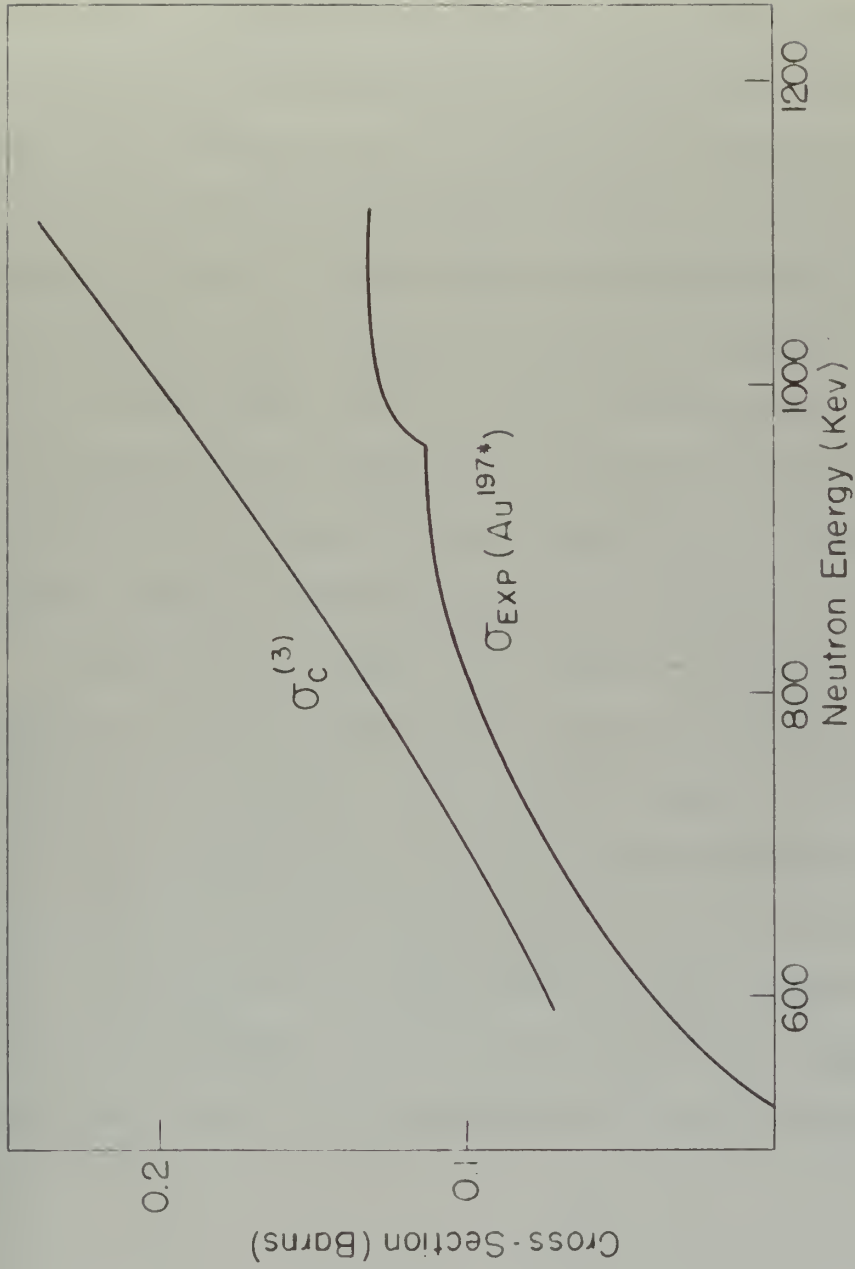


Figure 16

COMPARISON of EXPERIMENTAL CROSS-SECTION for PRODUCTION of Au^{197*} with $\sigma_c^{(3)}$ SHOWING THAT $\ell=3$ for the INCIDENT NEUTRONS IS POSSIBLE.

also shown on the same graph.

It is expected that the cross section for production of the isomeric state would be smaller compared to $\sigma_c^{(3)}$. However, the transition to the metastable state can compete very successfully with the transition to the ground state or other excited levels because it involves the emission of an $\ell = 0$ neutron. Decay of the compound nucleus to the ground state requires that an $\ell = 3$ neutron be emitted. Assuming that only the ground state and the metastable state compete, a calculation, such as described in Chapter II, reveals that 75 percent of the compound nuclei will go to the metastable level. Recalling that the experimental cross section may be in error by 50 percent in absolute magnitude, it is seen that the results are in agreement with the assumption that the activation takes place with an $\ell = 3$ neutron.

Figure 17 shows a plot of the neutron excitation curve from threshold to the first excited state above it, corrected for the cross section for the production of the compound nucleus. The resulting curve is the probability of emitting a neutron from the compound nucleus with $\ell = 0$ and the spin opposite to that of the incident neutron. This curve has the typical shape for the emission curve of an $\ell = 0$ neutron starting out as an $(E - E_{th})^{1/2}$ curve (also shown) and then falling below due to competition with decay to other levels.

also shown in the next graph.

It is assumed that the most probable the probability of the

infinite state (which is usually expressed by $\frac{1}{2}$). However, the

transition to the infinite state has a higher probability than

the transition to the finite state or other finite levels because it

involves the addition of an $\ell = 0$ neutron. (Many of the original

papers in the present state suggest that the $\ell = 1$ neutron be deleted.

Assuming that only the present state and the infinite state remain,

a calculation, such as described in Section II, reveals that the present

of the present state will be the infinite state. Assuming that

the equilibrium state is reached by the time the system is shown

this suggests, it is seen that the system has in equilibrium with the

assumption that the equilibrium state is reached with an $\ell = \frac{1}{2}$ neutron.

Figure 17 shows a plot of the neutron number versus time

relative to the time needed state shown is, compared for the same

method for the probability of the present neutron. The resulting

curve is the probability of finding a neutron from the original

neutron with $\ell = 0$ and the right appears to that of the infinite state

then. This curve has the typical shape for the solution curve of an

$\ell = 0$ neutron equation and as an $(1 - \frac{1}{2})^n$ curve (this shows) and

that falling below the $\ell = 0$ probability with time in other levels.

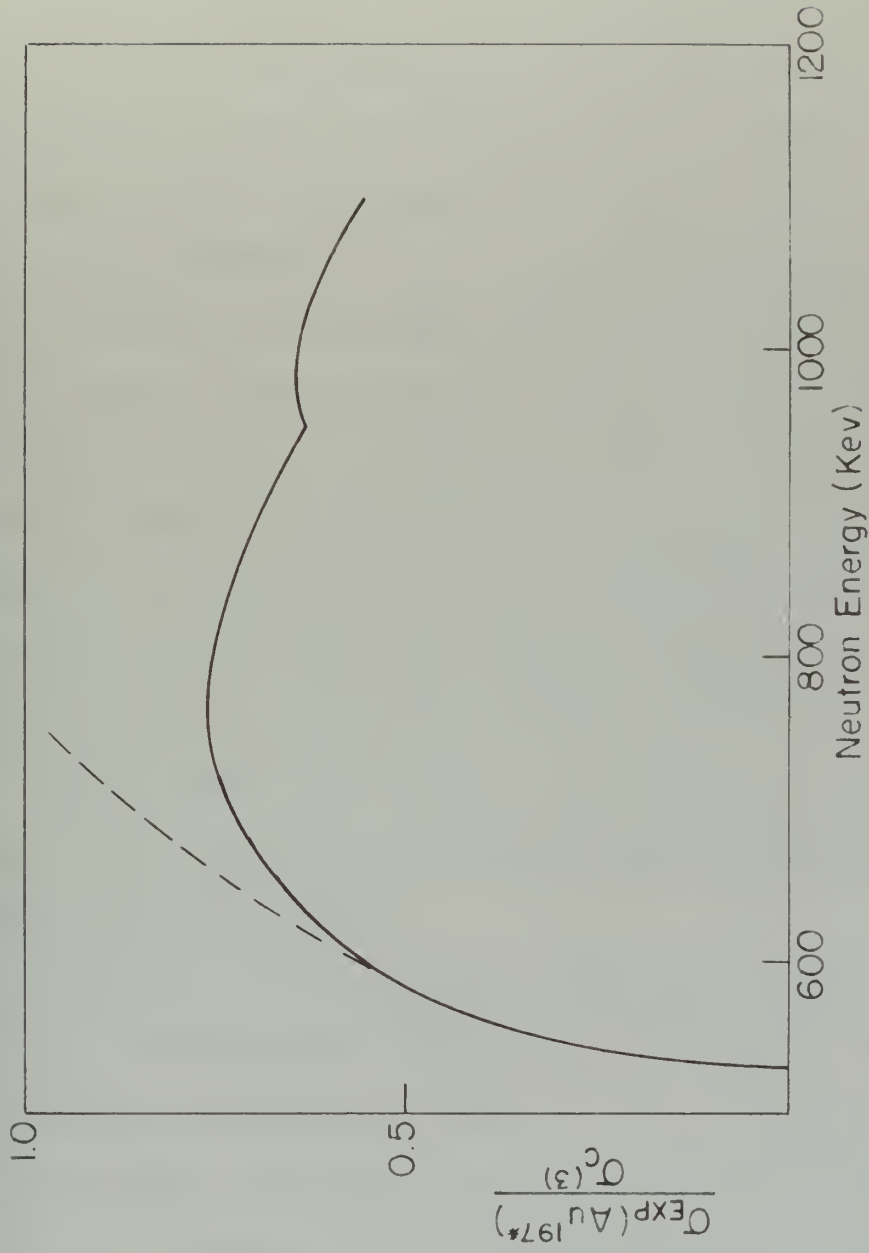


Figure 17

PROBABILITY OF COMPOUND NUCLEUS DECAYING TO METASTABLE STATE. THE BROKEN CURVE IS $Y = \text{const.} (E - E_{\text{TH}})^{1/2}$.

V. INVESTIGATION OF INDIUM

HISTORY OF In^{115*} EXCITATION

The metastable state in In^{115} has been studied by numerous investigators, not only to determine the properties of the state itself, but also as a convenient threshold detector. The relatively large cross section for production of the isomer and its convenient half-life make it ideal for many problems.

In 1939, Goldhaber, Hill, and Sillard⁴² reported the excitation of In^{115*} by inelastic scattering of neutrons from a radon-alpha-beryllium source giving neutrons of a maximum energy of about 13.5 Mev. The results were checked by exciting the state with 2.5-Mev d-d neutrons. It was shown that the activity belonged to the mass 115 indium isotope and not to the mass 113 by producing Cd^{115} from Cd^{116} by an (n,2n) reaction. The Cd^{115} they decayed to In^{115*} . On the other hand, Cd^{113} is stable and, hence, could not lead to In^{113*} .

The metastable state has since been produced by Barnes and Aradine⁴³ with 6.7-Mev protons and by Lark-Morovitz, Hasser, and Smith⁴⁴ with 16-Mev alpha-particles.

The first extensive excitation curve for this level was obtained by Waldman and Wiedenbeck⁴⁵. Using x-rays to produce the isomer, they reported levels at 1.12, 1.55, 2.13, and 2.64 Mev which decay to the metastable level. Subsequent measurements by Miller and Waldman⁴⁶ have confirmed the existence of the first two levels but corrected their energy to 1.04 and 1.42 Mev. The work by Waldman and Wiedenbeck re-

The word "state" is used in the sense of a political entity, not in the sense of a condition or state of mind. The word "state" is used in the sense of a political entity, not in the sense of a condition or state of mind.

In 1907, following the death of his father, he was elected to the position of Mayor of the City of New York.

1. The first group of people who were involved in the project were the members of the committee who were responsible for the selection of the project. They were the ones who decided that the project was worth doing and they were the ones who gave the project the go-ahead.

low levels of leachate must enter and while sufficient air

with 100% accuracy.

The first estimate was made by Mr. J. H. ...
by
... ..
... ..
... ..
... ..

quired the use of a thick x-ray target and therefore suffered from the same difficulty in interpretation as did Wiedenbeck's work on gold. However, Miller and Waldman found it possible to use a thin target that gave much more pronounced breaks in the excitation curve corresponding to the levels. In addition, they found a weak activity (twice the background) with 830-kev x-rays which was considered to be possibly the excitation of an 873-kev level, reported by Lawson and Cork⁴⁷.

A neutron excitation curve has been obtained by Cohen⁴⁸ using $C^{13}(d,n)$ and d-d neutrons from a 1-Mev Cavendish generator. Because of the reaction that was necessary to produce the neutrons, the curve could not be carried to the threshold of the excitation but had to be stopped at about 2 Mev. Figure 10 shows the results that were obtained. It is to be noted that the presence of individual levels was not detected. This is probably because so many levels contribute to the excitation at the high energy involved that the breaks due to the individual levels could not be seen.

In connection with an investigation of threshold detectors, Taschek⁴⁹ obtained four points on the neutron excitation curve in the energy region from 600 to 1500 kev. Of course, these data are too sketchy to show more than the general trend of the curve.

DESCRIPTION OF ENERGY LEVELS IN In^{115}

The energy of the isomeric state has been measured by Bell, Kettle, and Cassidy⁵⁰ using both a scintillation counter and a magnetic lens spectrometer. It is found that the K-conversion electrons have an energy of 312 kev, corresponding to an energy of 340 kev for the level.

[illegible]

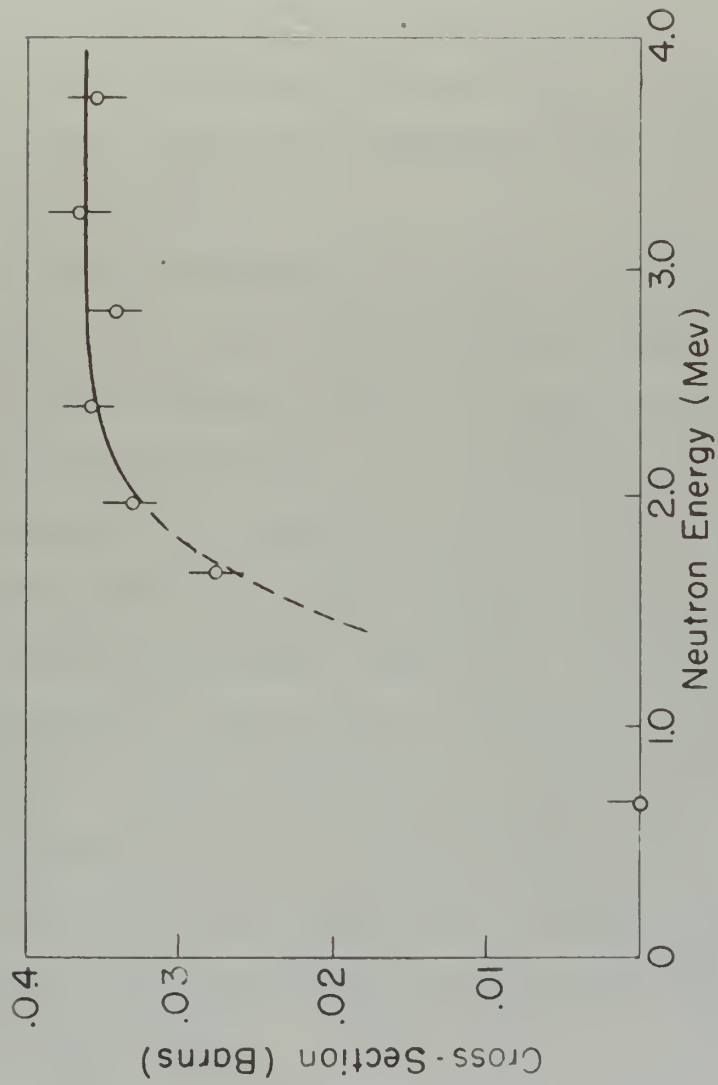


Figure 18

CROSS SECTION MEASURED BY COHEN for NEUTRON
EXCITATION of In^{115*} . (From Nature 161, 475 (1948))

The half-life of the state, according to the determination of Lawson and Cork⁵¹, is 4.5 hours.

Assuming the transition to be of the $\Delta = 4$ type, Δ defined in Chapter IV, the Weissacker hypothesis predicts a half-life of 4 seconds; while, if it is assumed to be $\Delta = 5$, the result is 300 hours. Although the agreement is not too good, it seems highly probable that the transition is of the latter type involving an electric 2^5 pole or a magnetic 2^4 pole radiation.

The K- to L-conversion ratio in this case does not give a conclusive decision between these two possibilities. Axel and Danneff⁵² report a theoretical value of 6.3 for the magnetic case and 2.64 for the electric multipole transition. The experimental value given is 5.0. Thus, it would seem to favor choosing the magnetic transition.

Consideration of the spins involved, however, removes the ambiguity. The ground state has a spin of $9/2$. The electric 2^5 pole transition requires a spin change of 5 between the ground and metastable states. If the possibility of a spin of $19/2$ is ruled out because it is too large, there is no possibility for an $\ell = 5$ transition. On the other hand, assuming the metastable level to have a spin of $1/2$ agrees with a magnetic 2^4 pole transition having a spin change of 4 and a parity change.

Figure 19 shows the metastable levels, together with those found at 600, 960, and 1370 kev from the neutron excitation curve determined in this investigation. Also shown is the competitive decay of the metastable level in about 5 percent of the cases by emission of a

The Δ - Δ transition is the only one which is not forbidden by the conservation of baryon number.

The following is a list of the names of the persons who have been
 elected to the office of the President of the United States since
 the year 1789. The names are given in the order in which they
 were elected.

The 1950-51 season was a very dry one and the water level in the reservoir was very low. The water level was only 1.5 feet above the normal level and the water was very muddy. The water level was only 1.5 feet above the normal level and the water was very muddy. The water level was only 1.5 feet above the normal level and the water was very muddy.

of 1 and a half dozen.

of 1/2 dozen with a separate 1/2 dozen (separate) being a 1/2 dozen
then. In the case of the separate 1/2 dozen, it was a 1/2
dozen in the large, there is no possibility of a 1/2 dozen -
table when. In the possibility of a 1/2 dozen in the large and
medium section, a 1/2 dozen of 1/2 dozen (separate) and 1/2 dozen
table. The 1/2 dozen has a 1/2 dozen of 1/2 dozen. The 1/2 dozen
possibilities of the 1/2 dozen, however, it was the 1/2 dozen

Figure 1 shows the estimated levels, together with lower and upper bounds, and 1970 law from the various countries were obtained in this investigation. This shows a low correlation of the estimated level in 1970 with the law of the country of a

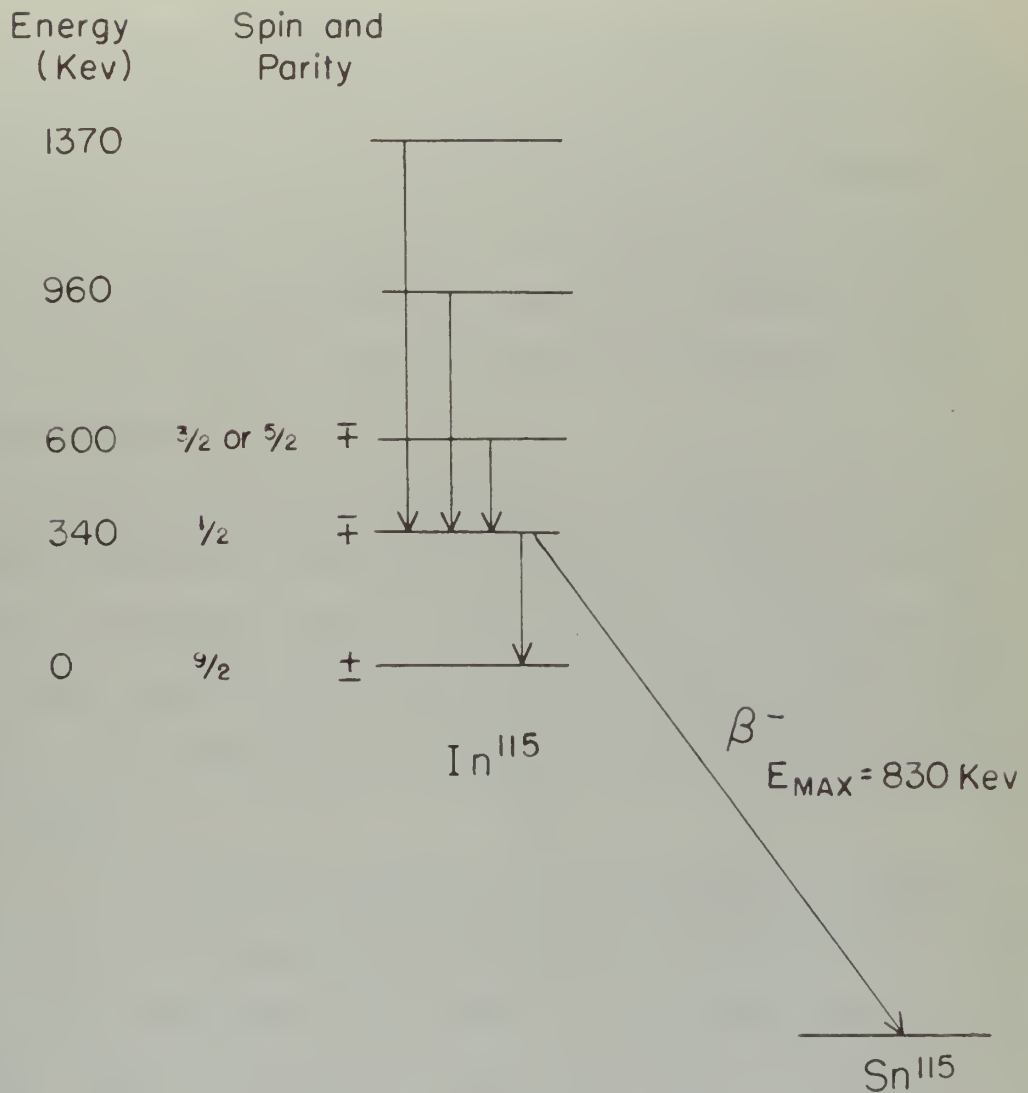


Figure 19

ENERGY LEVELS MEASURED in In¹¹⁵ by NEUTRON EXCITATION of the METASTABLE LEVEL.

beta-particle having a maximum energy of 830 kev, as found by Bell, Ketelle, and Cassidy⁵³.

The existence of this beta-decay has been confirmed in connection with this study of indium. Figure 20 shows a pulse-height distribution obtained from an 8-mil indium foil bombarded with neutrons from the M. I. T. cyclotron. The shape of the curve is distorted because the foil was thick compared with the range of the electrons, but the existence of the 312-kev conversion electron and 830-kev beta-decay is clearly shown.

It seems likely that the quantum-induced activity below 900 kev reported by Miller and Waldman and by Lawson and Cork was actually exciting of the 600-kev level. If this is assumed, it follows that the spin and parity of that level must be such that this process can take place often enough to be observed but that the process is somewhat forbidden.

A $\Delta = 2$ or lower transition (electric quadrupole or magnetic dipole) may be ruled out because this makes the transition to the ground state so much more probable than that to the metastable state that excitation would not be observed. On the other hand, $\Delta = 4$ or higher can be eliminated for the transition from the ground state to the 600-kev level because such a transition is too improbable for quantum excitation.

The remaining possibility $\Delta = 3$ involves a parity change and a spin change of 3 for electric 2^3 pole or of 2 for magnetic 2^2 pole

(Signature)

[illegible]

It seems likely that the possible solution to this problem is to be found in the fact that the present system of the law of the sea is based on the principle of the freedom of the seas. This principle is the basis of the law of the sea and is the basis of the law of the sea. It is the basis of the law of the sea and is the basis of the law of the sea.

low level because such a location is too important for the purpose
can be eliminated for the location from the present report to the div-
estimation would not be observed. In the other hand, $\Delta = \Delta$ is higher
again we need more evidence than that in the previous report that we
might say to reject and because this value the location for the present
 $\Delta = \Delta$ or lower location (whereas previous to reject)

The remaining possibility

$$K = \bigwedge_{i=1}^n p_i$$

also follows from the fact that K is a prime ideal.

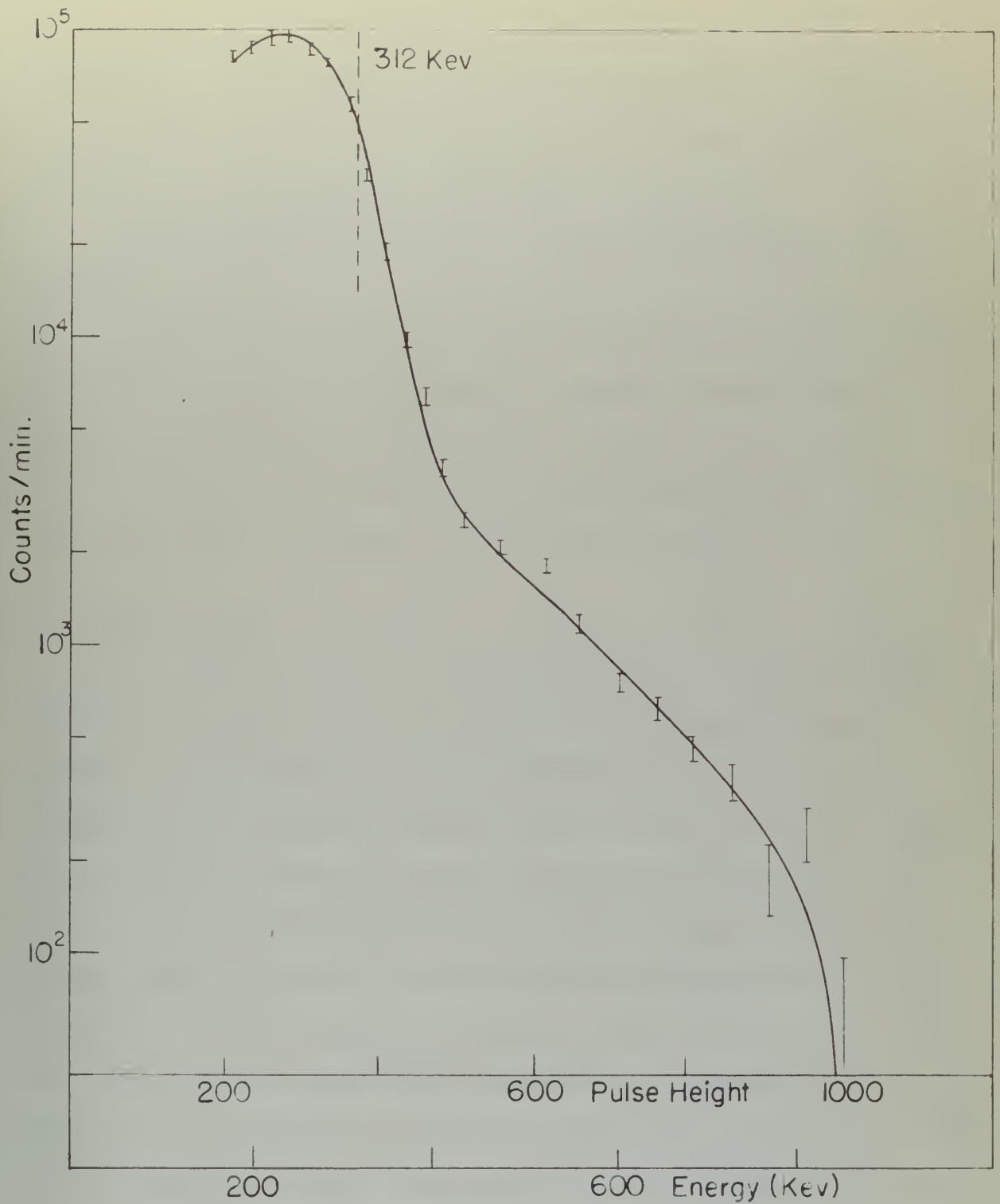


Figure 20

DIFFERENTIAL PULSE HEIGHT DISTRIBUTION for In^{115} from THICK FOIL SHOWING 312 Kev CONV. ELECTRON LINE and 830 Kev BETA DECAY (Energy calib. from Fig. 6)

excitation. A spin change of 2 with parity change could be provided by an $\ell = 1$ or an $\ell = 3$ neutron; if this is the case, neutrons will excite this level strongly. On the other hand, a spin change of 3 with parity change requires an $\ell = 3$ neutron, giving a small probability.

Unfortunately, it was not possible to determine the cross section for neutron excitation in this region. Therefore, no choice can be made between these two, and both spins are given on Figure 19.

For the case of the higher levels, many spin and parity choices will give results consistent with the observed excitation curves.

EXPERIMENTAL SETUP

The indium samples were in the form of foils 8 mils thick and 3 cm. diameter. These foils were irradiated for 4 hours with monoenergetic neutrons from the $\text{Li}^7(p,n)$ reaction using a 20-keV thick target on the Rockefeller generator. After irradiation, the foils were counted in the scintillation counter, described in Chapter III.

Such a long exposure, necessitated by the long lifetime of the isomer, introduces several problems into the determination of the excitation curve. It is impossible to keep the neutron flux constant over the exposure without sacrificing a portion of it to provide a range of adjustment. If the flux changes during this period, a simple integration of the neutron flux is not sufficient for the determination of the cross section. At the time these measurements were made, a continuously recording neutron-flux indicator was not available. It was therefore nec-

The first part of the paper is devoted to the study of the asymptotic behavior of the solutions of the system (1) as $t \rightarrow \infty$. It is shown that the solutions of the system (1) are bounded and tend to zero as $t \rightarrow \infty$. The second part of the paper is devoted to the study of the asymptotic behavior of the solutions of the system (1) as $t \rightarrow 0$. It is shown that the solutions of the system (1) are bounded and tend to zero as $t \rightarrow 0$.

will also provide protection with the colored warning marks.

For the case of the higher levels, even when the particle density was between 100 and 200 atoms per cubic cm (Figure 10).

The surface condition in this region, however, as shown was in conformity, it was not possible to determine the exact condition

[illegible]

essary to obtain several points on the excitation curve during the same irradiation, thus determining a section of the curve. These sections could then be joined by overlapping.

Another factor to be considered is the availability of generator time. If each point on the excitation curve is determined separately, the exposure time required becomes prohibitive.

For these reasons, advantage was taken of the dependence of the neutron energy on angle, and seven foils were exposed simultaneously at angles ranging from 0 to 75 degrees. Figure 21 shows the foil arrangement for exposure.

In order to obtain sufficient activity in the foils, it was necessary to place the foils so that the nearest edge was only 1 cm. from the neutron source. At this distance, the energy resolution was preserved by rolling the foils into small cylinders about 3 mm. in diameter and mounting them radially about the source.

The effect of capture of thermal neutrons was minimized by counting the side of the foil rolled inside. All of this capture activity is located in the very thin layer which was outside; so that most of the foil is shielded from thermal neutrons.

After irradiation the foils were counted in rotation six to eight times over a period of about 8 hours, starting with the lowest activity to separate the 54-minute capture activity from the 4.5-hour activity of In^{115*} . To increase the ratio of In^{115*} to In^{116} counts, advantage was taken of the energy sensitivity of the scintillation counter. The

[illegible]

...the ... of ...

The first meeting, which was held at the residence of the
 President, was held on the 1st of January, 1900, and was
 attended by the President, the Vice-President, the Secretary,
 and the members of the Executive Committee. The meeting
 was held in the morning, and the President presided.

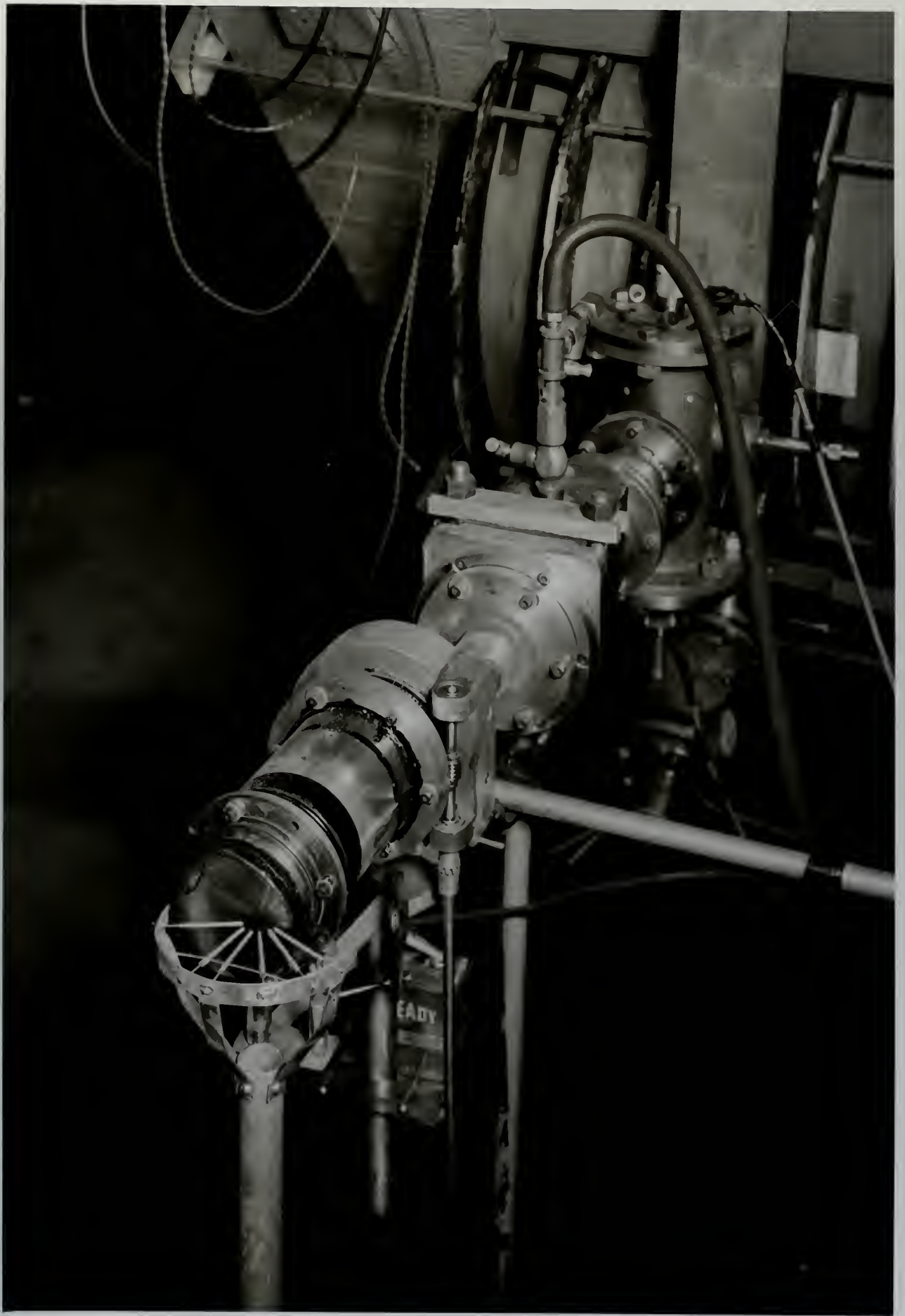
It seems to me that the only way to get the best results is to have the best people. I am not sure if this is the best way, but I think it is the only way. I am not sure if this is the best way, but I think it is the only way.

There is no evidence that the defendant was involved in the activities of the group. The defendant was not a member of the group and was not involved in the activities of the group. The defendant was not involved in the activities of the group.

[illegible]

Figure 21. Foil arrangement for the indium activation showing the foils rolled on the wire of the spider.

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DATE 08-19-2006 BY 60322 UCBAW



levels of the differential discriminator were set to record counts corresponding to an energy between 150 and 450 kev. This includes most of the counts from In^{115*} (see Figure 20) but includes only part of the counts from In^{116} ($E_{\text{max}} = 850$ kev for the beta-particle emitted).

An attempt to separate the two activities was made by comparing the number of counts above and below the pulse height corresponding to 450 kev, thus using the difference in shape of their pulse-height distributions. However, it was found that the statistical uncertainty in this method made it impracticable.

TREATMENT OF THE DATA

In order to take proper account of all points making up the decay curve of a given foil, it was considered desirable to compare the experimental curve with a theoretical decay curve. Since the construction of such a theoretical curve presupposes a knowledge of the relative amounts of the two activities present, a curve to include all cases had to be found.

All the experimental curves were plotted with the same scale on the same type of semilogarithm graph paper. A master decay curve was then constructed, also using the same scale, on a large sheet of paper starting with an assumption of the shorter-lived activity being 100 times that of the longer-lived and continued until the shorter-lived activity had decayed to 1/100 of the longer-lived. It is evident that, if the experimental curve is made up of a combination of the two activities within these limits, it must fit the master curve over some

level of the differential character was set to prevent further
 growing to an energy between 100 and 150 eV. This includes most of
 the energy from the ^{137}Cs (see Figure 10) but includes only half of the
 energy from the ^{137}Ba ($E_{\text{max}} = 500$ keV for the beta-decay method).
 In order to separate the two regions we used a computer
 the number of counts above and below the peak height corresponding to
 150 keV, then used the difference in terms of counts per channel
 distribution. However, it was found that the statistical uncertainty
 in this method was too large.

TREATMENT OF THE DATA

In order to make proper account of all points within the energy
 range of a given peak, it was necessary to compare the exper-
 imental curve with a theoretical decay curve. Since the theoretical
 at each a theoretical curve represents a knowledge of the relative
 amounts of the two activities present, a curve is chosen all cases had
 to be found.

All the experimental curves were plotted with the same scale on
 the same type of logarithmic graph paper. A series of curves were
 then constructed, also using the same scale, on a large sheet of paper
 starting with an assumption of the number of activities being 100
 that that of the longer-lived and continued until the shorter-lived
 activity had decayed to 1/100 of the longer-lived. It is evident that
 if the experimental curve is made up of a combination of the two activ-
 ities which decay faster, it must fit the shorter curve very well.

interval. By superposing the experimental curve on the master and by sliding the former to find the best fit, each point could be given its proper weight. To permit determination of the amount of each activity, two lines were drawn on the master corresponding to the contribution of each activity to the total curve. It was then a simple matter to read off the amount of each activity at any time and in particular at the time corresponding to $t = 0$ on the experimental plot.

In this investigation, the 105-minute isomeric activity in In^{113} was neglected. That this would be possible was predicted from the relative abundance (4.2 percent for In^{113} , compared with 95.8 percent for In^{115}). The decay curves obtained justify this assumption, inasmuch as the combination of the 54-minute and the 4.5-hour activities were sufficient in all cases to explain their shape.

The accuracy of the method described above was high enough so that the uncertainty in the determination of the In^{115*} activity was estimated to be less than 10 percent.

Since the neutron flux is anisotropic, the various activities must be corrected as a function of the exposure angle. Taschek and Homendinger⁵⁴ have shown that the flux is not even isotropic in the center-of-mass coordinates for the $\text{Li}^7(p,n)$ reaction because of the resonance in the compound nucleus. The flux as a function of angle must therefore be measured. Taschek and Homendinger's measurements do not extend into the region of interest in this work. In conjunction with Willard, the neutron flux as a function of angle from 0 to 90

hence, the expression for the potential energy in the region V is

$$U = \frac{1}{2} \int_V \epsilon |\mathbf{E}|^2 dV + \frac{1}{2} \int_V \mu |\mathbf{H}|^2 dV + \frac{1}{2} \int_V \rho \phi dV + \frac{1}{2} \int_V \sigma \psi dV$$
 where ϕ and ψ are the electrostatic and magnetostatic potentials respectively. The first two terms are the energy of the electric and magnetic fields respectively. The last two terms are the energy of the charges and currents respectively. The total energy is then given by

$$U = \frac{1}{2} \int_V \epsilon |\mathbf{E}|^2 dV + \frac{1}{2} \int_V \mu |\mathbf{H}|^2 dV + \frac{1}{2} \int_V \rho \phi dV + \frac{1}{2} \int_V \sigma \psi dV$$
 This expression is valid for any region V in which the fields are defined.

In this section, we shall consider the case in which the fields are defined in the region V and the charges and currents are distributed in the region V . The total energy is then given by

$$U = \frac{1}{2} \int_V \epsilon |\mathbf{E}|^2 dV + \frac{1}{2} \int_V \mu |\mathbf{H}|^2 dV + \frac{1}{2} \int_V \rho \phi dV + \frac{1}{2} \int_V \sigma \psi dV$$
 The first two terms are the energy of the electric and magnetic fields respectively. The last two terms are the energy of the charges and currents respectively. The total energy is then given by

$$U = \frac{1}{2} \int_V \epsilon |\mathbf{E}|^2 dV + \frac{1}{2} \int_V \mu |\mathbf{H}|^2 dV + \frac{1}{2} \int_V \rho \phi dV + \frac{1}{2} \int_V \sigma \psi dV$$
 This expression is valid for any region V in which the fields are defined.

The energy of the system is then given by

$$U = \frac{1}{2} \int_V \epsilon |\mathbf{E}|^2 dV + \frac{1}{2} \int_V \mu |\mathbf{H}|^2 dV + \frac{1}{2} \int_V \rho \phi dV + \frac{1}{2} \int_V \sigma \psi dV$$
 The first two terms are the energy of the electric and magnetic fields respectively. The last two terms are the energy of the charges and currents respectively. The total energy is then given by

$$U = \frac{1}{2} \int_V \epsilon |\mathbf{E}|^2 dV + \frac{1}{2} \int_V \mu |\mathbf{H}|^2 dV + \frac{1}{2} \int_V \rho \phi dV + \frac{1}{2} \int_V \sigma \psi dV$$
 This expression is valid for any region V in which the fields are defined.

Thus, the energy of the system is given by

$$U = \frac{1}{2} \int_V \epsilon |\mathbf{E}|^2 dV + \frac{1}{2} \int_V \mu |\mathbf{H}|^2 dV + \frac{1}{2} \int_V \rho \phi dV + \frac{1}{2} \int_V \sigma \psi dV$$
 The first two terms are the energy of the electric and magnetic fields respectively. The last two terms are the energy of the charges and currents respectively. The total energy is then given by

$$U = \frac{1}{2} \int_V \epsilon |\mathbf{E}|^2 dV + \frac{1}{2} \int_V \mu |\mathbf{H}|^2 dV + \frac{1}{2} \int_V \rho \phi dV + \frac{1}{2} \int_V \sigma \psi dV$$
 This expression is valid for any region V in which the fields are defined.

degrees was measured for proton energies from the $\text{Li}^7(p,n)$ threshold to 3.75 Mev using the long counter described in Chapter III. The scattered neutron background was determined by use of a paraffin cone to shadow the counter. The results of these measurements are reported by Willard⁵⁵, and are plotted here in Figure 21a.

Each irradiation of foils covered an energy interval of from 300 to 500 keV. The intervals were chosen so that they overlapped approximately half of their extent. Each segment of the In^{115*} excitation curve determined in the interval was corrected for the anisotropy of the flux and energy resolution and was then normalized to the preceding segment. This technique involves a propagation of errors but does not affect the location of energy levels found by means of the curve. Figure 22 shows the complete curve.

DISCUSSION OF THE RESULTS

In spite of the scattering of the points on the excitation curve (Figure 22), the presence of three levels is shown at 600 keV, 960 keV, and 1.37 MeV with possibly a fourth level at 1.75 MeV. The uncertainty in these energy assignments is estimated to be of the order of 40 keV. This uncertainty, however, is considered hardly sufficient to bring the 960 keV and the 1.37 MeV level into agreement with Miller and Waldman's values of 1.02 and 1.42 MeV for quantum-excited levels. The correspondence is close enough to leave little doubt that the same levels are being excited and suggests that the energy calibration for one of the experiments is in error. Unfortunately, there is not a third accurate

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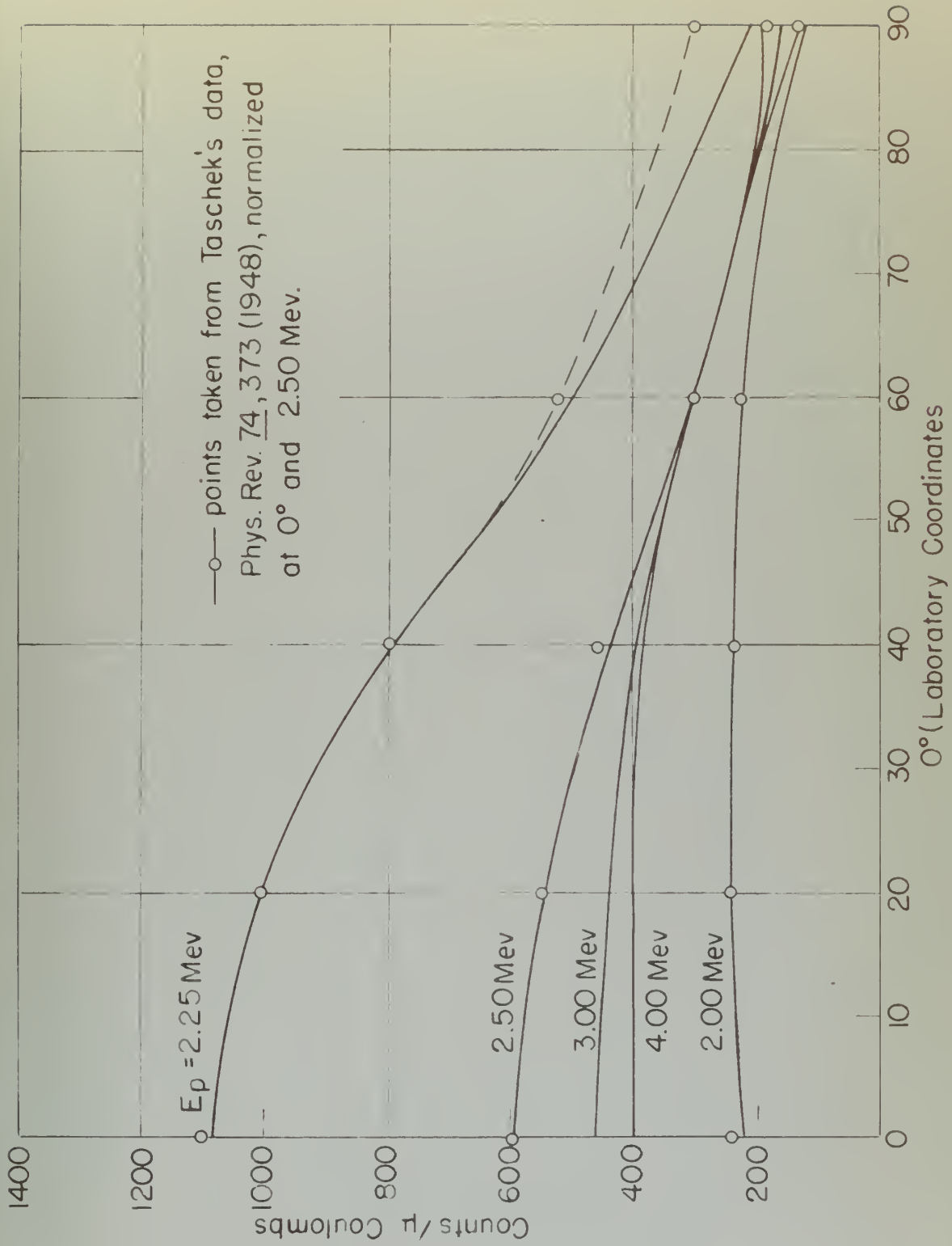


Figure 21a

RELATIVE DIFFERENTIAL CROSS SECTION for the $\text{Li}^7(\text{p},\text{n})\text{Be}^7$ REACTION

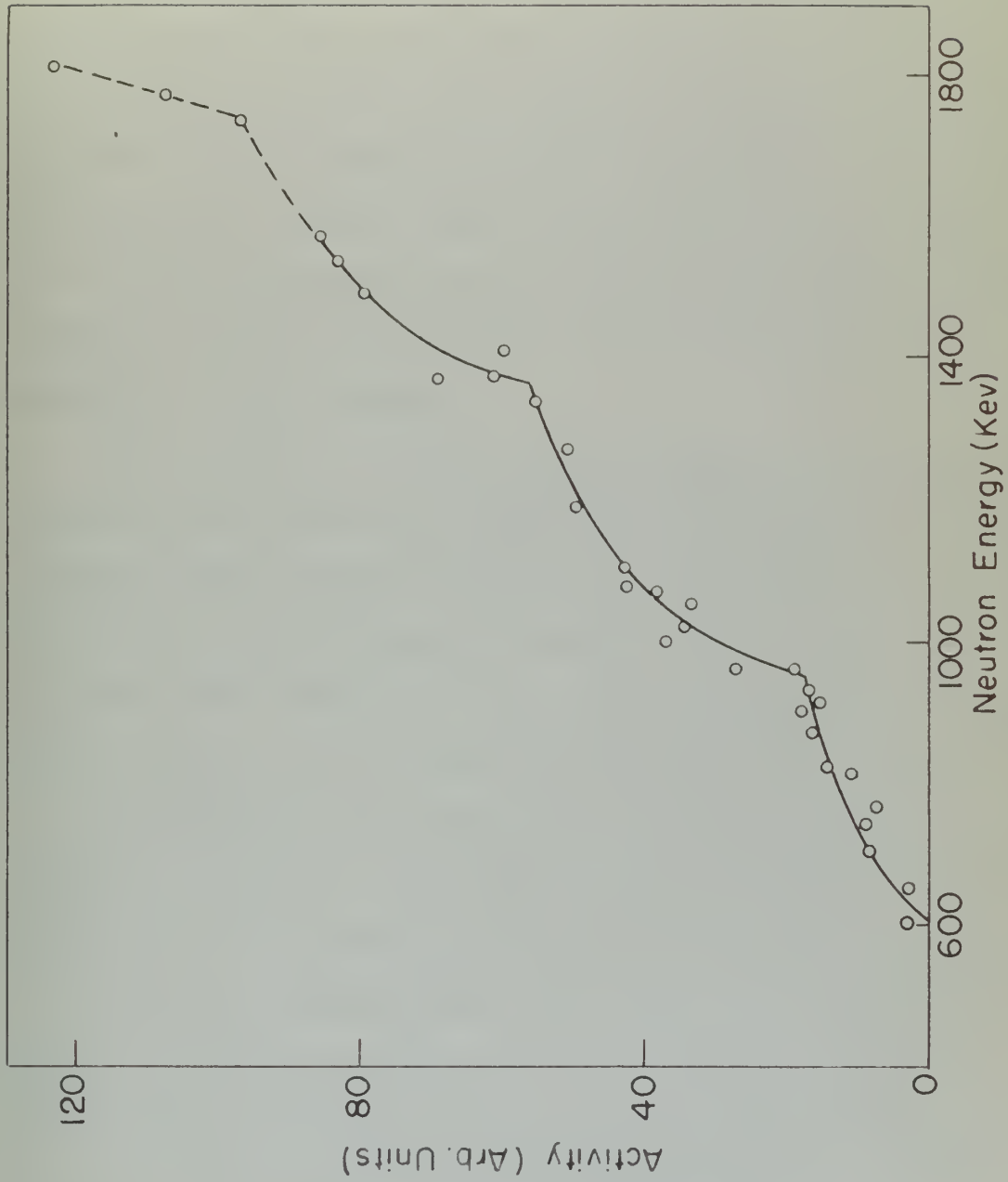


Figure 22
NEUTRON EXCITATION CURVE for In^{115*}

determination of either of these levels for comparison. From the reproducibility of the calibration of the Rockefeller generator, it is difficult to see how it could be in error by the amount suggested above.

Probably the greatest source of error in the individual points on the excitation curve is due to the stringent requirements on positioning the foils. Since the point at which the foil may be considered concentrated for the irradiation is only 2.1 cm. from the target (found by integrating an inverse square activity over a circular foil), a 2-mm. error in the mounting of the foils or in determining the point at which the proton beam strikes the target results in a 20 percent error in the induced activity.

The location of the beam cannot be observed after the target has been installed so that it is impossible to know whether the beam has wandered slightly during the irradiation. This error can be corrected for, however, because foils were exposed at angles on both sides of 0 degrees. It was then assumed that the beam actually hit the target at such a point that the activities were symmetrical about 0 degrees.

The placement of the individual foils on the circumference of a circle with its center at the target could not be checked in such a simple fashion and is no doubt the reason for the spread in the points.

In further investigation of this problem, it is suggested that a thicker target (of the order of 50 kev) be used and the foils be mounted farther from the target, even at the expense of reducing the activity below that obtained in this investigation.

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VI. CONCLUSIONS AND SUGGESTIONS FOR FURTHER INVESTIGATION

SUMMARY OF RESULTS

Neutron excitation reveals the existence of energy levels in Au^{197} at 1.14 Mev and 1.14 Mev. It is found that the metastable state at 540 kev can be excited directly by $\ell = 3$ neutrons. The spin of this level is shown to be $11/2$ with parity opposite that of the ground state. This fixes the spin of the 270-kev level to which the metastable level decays at $5/2$ with parity also opposite that of the ground state.

Since the 1.14-Mev level is also excited by photons, its spin is shown to be $7/2$ or $9/2$ with parity opposite to that of the ground state.

An analysis of the data for direct production of the metastable state shows the cross section to be reasonably in agreement with that predicted by the continuum theory for the production of the compound nucleus by an $\ell = 3$ neutron.

The levels found in In^{115} by neutron excitation are at 600 kev, 960 kev, and 1370 kev. The latter two are also found by quantum excitation, but the quantum values are slightly higher (8 percent). An analysis shows that the metastable level has a spin equal to $1/2$, and the 600-kev level, a spin of $5/2$ or $3/2$. Both of these levels have a parity different from the ground state.

FURTHER INVESTIGATION OF Au^{197}

Because of the disagreement between energy levels obtained in the investigation and those obtained by Wiedenbeck using thick-target x-rays,

THE HISTORY OF THE UNITED STATES

CHAPTER I

The first of the great principles of the American Revolution was the right of the people to alter or to abolish their government, and to institute a new one, when it became necessary for them to do so. This principle was the basis of the Declaration of Independence, and it was the first step towards the establishment of a new and better government for the United States.

The second principle was the right of the people to be represented in their government. This principle was the basis of the Constitution, and it was the second step towards the establishment of a new and better government for the United States.

CHAPTER II

The third principle was the right of the people to be governed by laws which they had a part in making. This principle was the basis of the Constitution, and it was the third step towards the establishment of a new and better government for the United States.

it would seem of great value to redetermine the quantum excitation curve using x-rays from a thin target. Determination of the relative probability for exciting the various levels in this fashion would indicate possible spin and parity values for some of the higher excited levels.

An extension of the neutron excitation curve to higher energies would probably reveal the presence of other levels. A logical choice for the neutron source in such a study is the d-d reaction which will yield from 2.5-Mev to 7-Mev neutrons with the Rockefeller generator. The use of deuterons as the particle accelerated has the disadvantage of giving very high backgrounds of neutrons even though the beam is not striking the target. This is partly because so many of the deuteron reactions have low thresholds. In addition, deuterons in the beam are imbedded in the metal a distance corresponding to their range. When the accelerating voltage is raised slightly, this forms an effective target for the d-d reaction.

However, if a counter for the gold activity can be found that is insensitive to this background or if it is possible to cut off the beam farther from the counter, it will be possible to extend the results of this investigation.

FURTHER INVESTIGATION OF In^{115*}

Better results could be obtained for the excitation of In^{115*} if the positioning error could be reduced by exposing the foils at a greater distance from the target. In Chapter V, it was suggested that

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this be accomplished by use of a thicker lithium target. In addition, the beam current available from the Rockefeller generator has been increased by a factor of 3 or 4 since the exposure were made. It is felt that, with a 50-kev lithium target and the present beam current on the order of 5 microamperes, sufficient activity could be obtained in the foils when they are mounted with the nearest edge on the order of 6 cm. from the target.

By connecting the output of the long counter to a counting-rate meter that operates a continuously recording meter, such as an Esterline-Angus pen recorder, the neutron flux on the foils would be determined as a function of time. This, plus a measurement of the efficiency of the In^{115*} activity counter, would enable a determination to be made of the cross section for neutron excitation of the state as a function of energy. A comparison of this cross section with the theoretical value will remove the ambiguity in the spin of the 600-kev level.

INVESTIGATION OF OTHER ISOTOPES

Several other isotopes (see Table I) lend themselves to an investigation similar to that made on gold and indium. All may either be treated according to the technique used for gold or that used for indium.

Rhodium has the advantage of being composed of 99.9 percent of the isotope containing the metastable state, Rh^{103} . The excitation by x-rays has been accomplished by Wiedenbeck, as reported in Chapter I. The metastable state is only of the order of 60 kev above the ground state, but radiation of this energy can be detected easily by a scintil-

There is no doubt that the above is a very good example of the work of a good writer. The style is clear and concise, and the facts are presented in a logical and easy-to-understand manner. The use of simple language and short sentences makes the text accessible to a wide range of readers. The overall impression is one of a well-written and informative piece of work.

It is a very common mistake to suppose that the only way to get the best of the law is to get the best of the facts. This is not so. The law is a system of rules, and the facts are the material upon which the law is applied. The law is a system of rules, and the facts are the material upon which the law is applied. The law is a system of rules, and the facts are the material upon which the law is applied.

The following is a list of the names of the persons who have been identified as having been in contact with the subject of this report, in the period from 1964 to 1968. The names are listed in alphabetical order, and are followed by the date of the last contact with the subject of this report.

lation counter. Its half-life of 45 minutes requires that any extensive investigation be carried on in a fashion similar to that used for indium.

The excitation of Cd^{111} by Wiedenbeck has also been reported in Chapter I. This state has an energy of 14.9 kev with a half-life of 40.7 minutes. In addition, a metastable state in Cd^{113} has been found by Helmholtz and McGinnis⁵⁶ with a half-life of 2.3 minutes. The cross sections for both of these activities are expected to be small because they occur in isotopes with only about 12 percent relative abundance. However, there is no activity expected that would confuse the results, and their lifetimes are such that counting may be made at some distance from the generator, resulting in low background.

Another possible choice for investigation is Ba^{137} with a 156-second half-life. Unfortunately, this isotope is also low in relative abundance, being 11 percent. There have been no investigations of the excitation curve reported for this state.

Hafnium has a 19-second, 200-kev activity assigned to either Hf^{177} or Hf^{179} . The former is 13 percent abundant and the latter, 14 percent. Although it may prove possible to lower the background during the counting period by removing the accelerator voltage, it is more likely that the beam will have to be cut off by a flap valve as in the case of gold. If this is true, the discussion for wolfram given below also applies to hafnium.

[illegible]

Last among the logical choices is W^{183} . This isotope of wolfram could be investigated with the technique used for gold, since its half-life is 5.5 seconds. The isotopic abundance of 14 percent will result in a low cross section in all probability. Unlike barium and cadmium, however, the lifetime is so short that the voltage cannot be removed from the accelerator, and backgrounds on the order of those obtained for gold will be encountered. If the same cross section per nucleus is assumed for W^{183} as was found for Au^{197} , the activity expected is so low that the determination of the wolfram curve is a marginal case.

In summary, if several rhodium foils could be obtained, Rh^{103} would seem the best choice for future investigation, followed by the two cadmium isotopes and Ba^{137} . Hafnium and wolfram appear to pose a very difficult problem with the present neutron source strength.

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VII. ACKNOWLEDGMENTS

I wish to thank Dr. Clark Goodman for his advice, interest, and constructive criticism during the progress of this research and the preparation of this thesis. Drs. Victor F. Weisskopf and Harvin L. Goldberger contributed many valuable suggestions on the theoretical aspects of the problem. I am also indebted to Dr. William W. Buechner for helpful discussions about this research.

These experiments could not have been performed without the generous support of the Nuclear Shielding Project at M. I. T. under the sponsorship of the U. S. Navy, Bureau of Ships, and Office of Naval Research.

My thanks are due Dr. Truman S. Gray and the group working with him for their assistance in assembling the equipment. Mr. H. B. Frey provided invaluable aid in the building of the scintillation spectrometer and in the electronic problems with the associated apparatus.

Mr. Harvey B. Willard assisted in many of these measurements and contributed many worthwhile suggestions.

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I thank Dr. William M. Preston and the men who work with him, Messrs. Donald C. Thompson, Ira M. Slawson, Vincent Yaras, and Edgar Jansen, for their assistance in operation of the Rockefeller generator. I also thank Dr. David H. Frisch for the rotating target that made possible the neutron intensity that was needed.

I wish to thank Mr. James G. Thompson for his interest, assistance and the cooperation of this branch. The following is a list of the names of the persons who have been in contact with the branch since the beginning of the year.

These accounts could not have been written without the
generous support of the Society. I am, therefore, indebted to it for the
the opportunity of the U. S. Navy, Bureau of Ships, and Office of
Naval Research.

It should be noted that the above information is based on the information provided by the individual named in the caption of this document. The information is not intended to be used for any other purpose.

10. The above information is true and correct to the best of my knowledge and belief.

the following information:

I thank Mr. William H. Taft for his kind and able help.
Very truly yours,
John D. Thompson, Jr.
His Excellency, William H. Taft, Secretary of the Interior,
Washington, D. C.

My appreciation is due Mrs. Mary E. White and Mrs. Charles Rowe, Jr. for their excellent work in the preparation of the manuscript.

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I should also like to express my indebtedness to Dr. Nathaniel H. Frank and Dr. George G. Harvey for their guidance and suggestions throughout my graduate work at M. I. T.

It is requested that the Board of Directors of the
Company be authorized to execute and deliver such
instruments as may be necessary to carry out the
above purposes.

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1. The first part of the book is devoted to a general introduction to the subject of the history of the English language. It discusses the various factors which have influenced the development of the language, such as contact with other languages, internal changes, and the influence of social and cultural factors. The author also discusses the importance of the study of the history of the English language for the understanding of the language itself.
2. The second part of the book is devoted to a detailed study of the history of the English language from the beginning of the 15th century to the present. It discusses the various stages of the language, from Old English to Middle English to Modern English, and the changes which have taken place in the language over time. The author also discusses the influence of other languages on the English language, such as Latin, French, and Greek.
3. The third part of the book is devoted to a study of the history of the English language in the United States. It discusses the various factors which have influenced the development of the American English language, such as contact with other languages, internal changes, and the influence of social and cultural factors. The author also discusses the importance of the study of the history of the American English language for the understanding of the American English language itself.
4. The fourth part of the book is devoted to a study of the history of the English language in the British Isles. It discusses the various factors which have influenced the development of the British English language, such as contact with other languages, internal changes, and the influence of social and cultural factors. The author also discusses the importance of the study of the history of the British English language for the understanding of the British English language itself.
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6. The sixth part of the book is devoted to a study of the history of the English language in the world. It discusses the various factors which have influenced the development of the world English language, such as contact with other languages, internal changes, and the influence of social and cultural factors. The author also discusses the importance of the study of the history of the world English language for the understanding of the world English language itself.
7. The seventh part of the book is devoted to a study of the history of the English language in the future. It discusses the various factors which will influence the development of the future English language, such as contact with other languages, internal changes, and the influence of social and cultural factors. The author also discusses the importance of the study of the history of the future English language for the understanding of the future English language itself.
8. The eighth part of the book is devoted to a study of the history of the English language in the present. It discusses the various factors which have influenced the development of the present English language, such as contact with other languages, internal changes, and the influence of social and cultural factors. The author also discusses the importance of the study of the history of the present English language for the understanding of the present English language itself.
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1. The first part of the document is a list of names and dates, arranged in two columns. The names are written in a cursive script, and the dates are in a more formal, printed style. The names are: John Smith, James Brown, William Jones, Thomas White, Robert Black, and Charles Green. The dates are: 1789, 1790, 1791, 1792, 1793, and 1794.
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7. The seventh part of the document is a list of names and dates, arranged in two columns. The names are written in a cursive script, and the dates are in a more formal, printed style. The names are: John Smith, James Brown, William Jones, Thomas White, Robert Black, and Charles Green. The dates are: 1789, 1790, 1791, 1792, 1793, and 1794.
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9. The ninth part of the document is a list of names and dates, arranged in two columns. The names are written in a cursive script, and the dates are in a more formal, printed style. The names are: John Smith, James Brown, William Jones, Thomas White, Robert Black, and Charles Green. The dates are: 1789, 1790, 1791, 1792, 1793, and 1794.
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1. The first part of the report deals with the general situation of the country and the progress of the work during the year.

2. The second part contains a detailed account of the work done in the various departments of the institution.

3. The third part gives a summary of the results of the work and a statement of the financial position of the institution.

4. The fourth part contains a list of the names of the persons who have been employed by the institution during the year.

5. The fifth part contains a list of the names of the persons who have been elected to the various offices of the institution.

6. The sixth part contains a list of the names of the persons who have been elected to the various committees of the institution.

7. The seventh part contains a list of the names of the persons who have been elected to the various boards of the institution.

8. The eighth part contains a list of the names of the persons who have been elected to the various councils of the institution.

9. The ninth part contains a list of the names of the persons who have been elected to the various assemblies of the institution.

10. The tenth part contains a list of the names of the persons who have been elected to the various synods of the institution.

11. The eleventh part contains a list of the names of the persons who have been elected to the various conferences of the institution.

12. The twelfth part contains a list of the names of the persons who have been elected to the various conventions of the institution.

13. The thirteenth part contains a list of the names of the persons who have been elected to the various assemblies of the institution.

14. The fourteenth part contains a list of the names of the persons who have been elected to the various synods of the institution.

15. The fifteenth part contains a list of the names of the persons who have been elected to the various conferences of the institution.

16. The sixteenth part contains a list of the names of the persons who have been elected to the various conventions of the institution.

17. The seventeenth part contains a list of the names of the persons who have been elected to the various assemblies of the institution.

18. The eighteenth part contains a list of the names of the persons who have been elected to the various synods of the institution.

19. The nineteenth part contains a list of the names of the persons who have been elected to the various conferences of the institution.

20. The twentieth part contains a list of the names of the persons who have been elected to the various conventions of the institution.

BIOGRAPHICAL SKETCH

Born: Waterloo, Iowa, June 11, 1920

Education:

West High School, Waterloo, Iowa, graduated 1938

Iowa State Teachers College, Cedar Falls, Iowa, B. A., 1942

Naval Post-Graduate School, Annapolis, Maryland, 1946-1947

Honors:

Tuition Scholarship, Iowa State Teachers College, 1941-1942

Professional Experience:

United States Navy 1942 -

Present Rank of Lieutenant

Associations:

Associate Member, The Society of the Sigma Xi

Member, The American Physical Society

STUDY OF THE PROBLEM

Abstract, 1951, 1952, 1953

1951

1952

1953

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